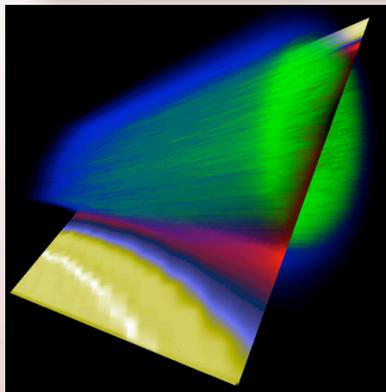


This research used resources of the Argonne Leadership Computing Facility at Argonne National Laboratory, which is supported by the Office of Science of the U.S. Department of Energy under contract DE-AC02-06CH11357.

LLNL-PRES-558099

NIC

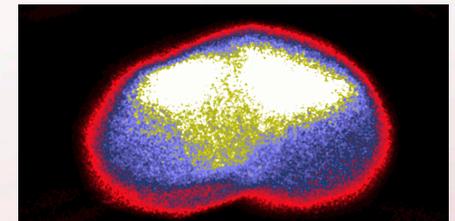
Laser-Plasma Interactions (LPI) in National Ignition Campaign (NIC) Hohlräume



D. E. Hinkel

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AX Division

Weapons and Complex Integration Directorate
Lawrence Livermore National Laboratory



In collaboration with:

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C. H. Still, D. A. Callahan, L. Divol, S. H. Langer, L. J. Suter**

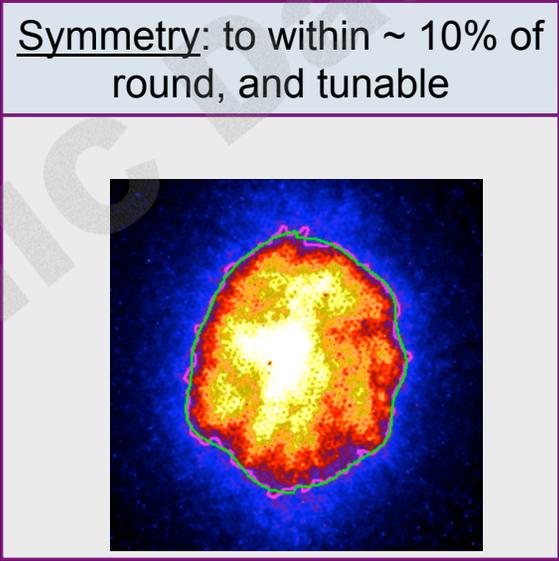
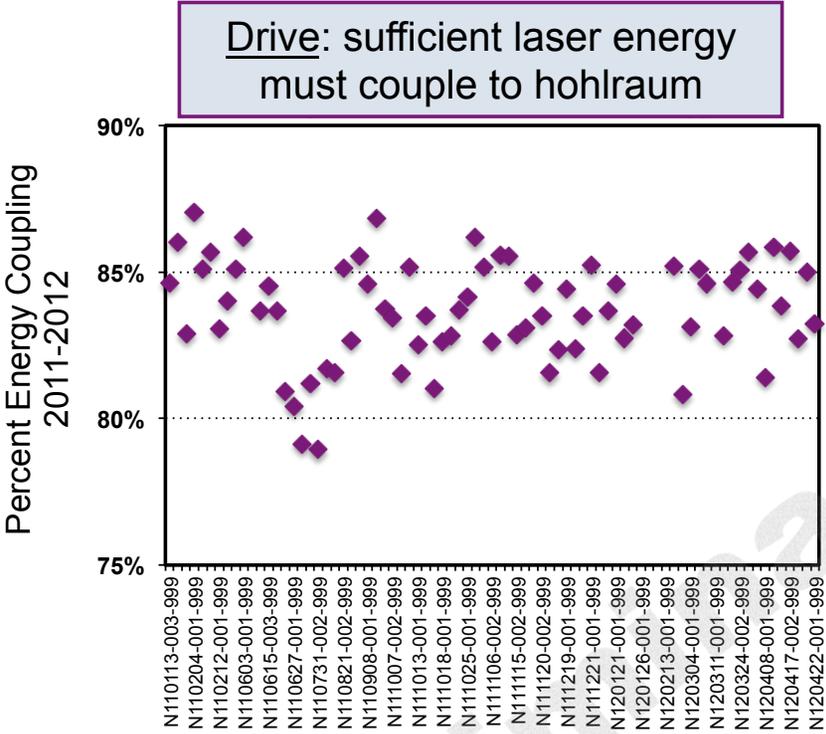
May 10, 2012

Lawrence Livermore National Laboratory · National Ignition Campaign

This work performed under the auspices of the U.S. Department of Energy by Lawrence Livermore National Laboratory under Contract DE-AC52-07NA27344

The NIC hohlraum provides high radiation drive and near-round capsule implosions

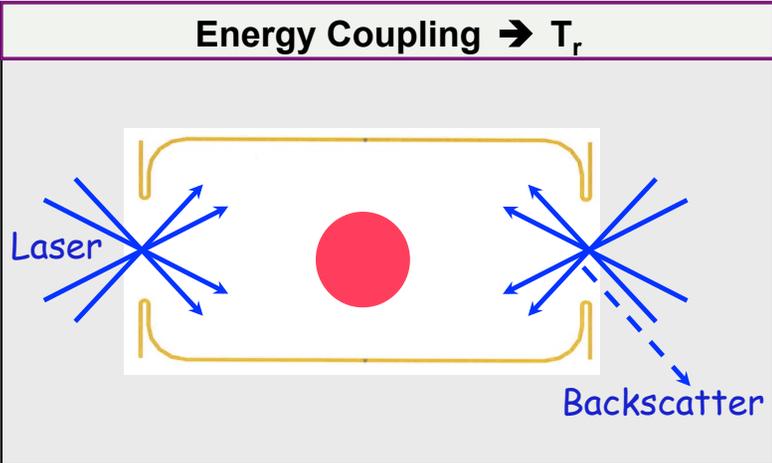
- To meet drive and symmetry requirements:



- NIC Laser-Plasma Interactions (LPI): reduce energy coupling; provide symmetry tuning

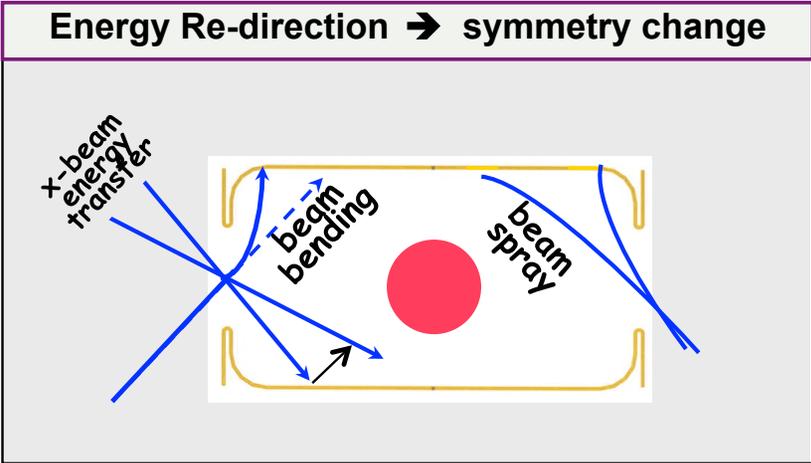
Can we improve coupling?

LPI can modify how laser energy couples to the hohlraum target



SBS: laser scatters off self-generated ion acoustic waves (iaws)

SRS: laser scatters off self-generated electron plasma waves (epws)



Beam spray: intensifies & scatters light

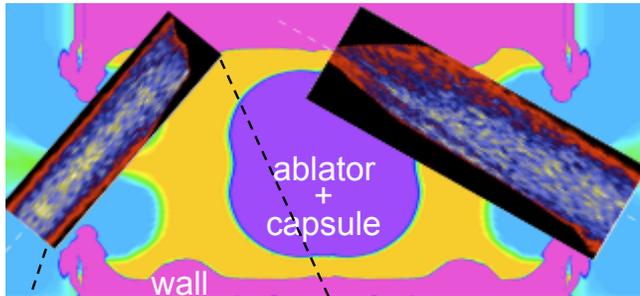
Beam bending: moves light pattern on wall

Cross-beam transfer: moves laser power among beams

Re-absorption of scattered light: alters symmetry

LPI processes span a wide range of length and time scales

macroscale

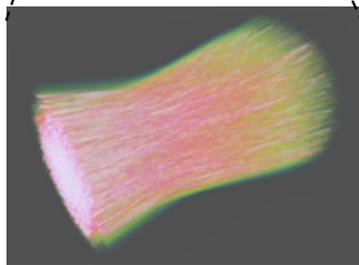


- Hydrodynamic length and time scales are set by target size [$O(\text{mm})$] and laser pulse length [$O(\text{ns})$]

environment -- plasma parameters and scale lengths

need accurate plasma conditions

mesoscale

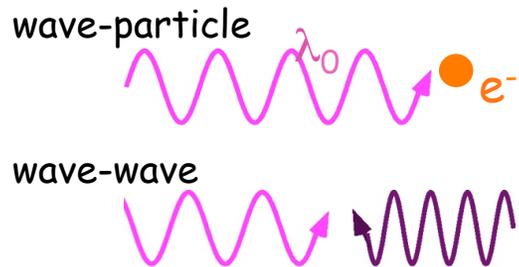


- LPI evolves on: μm length and ps time scales

beam propagation

need accurate beam model (including effects of cross-beam energy transfer, pump depletion, competition)

microscale



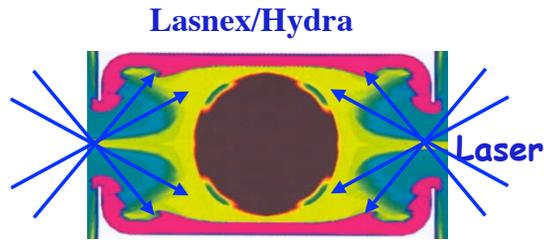
- Detailed processes of LPI occur on “light” spatial and temporal scales

kinetic effects

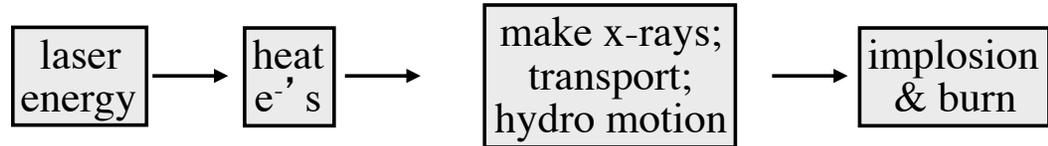
need accurate interactions – when do kinetic effects matter?

Our approach to multi-scale modeling uses a suite of tools

macroscale

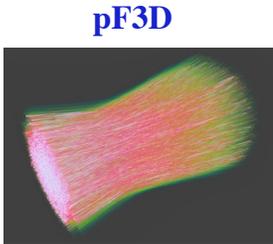


Radiation/Hydrodynamics simulations

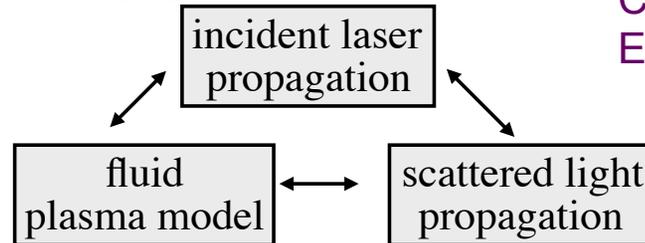


Calculate Gain Exponent G : $R = \eta \exp(G)$

mesoscale

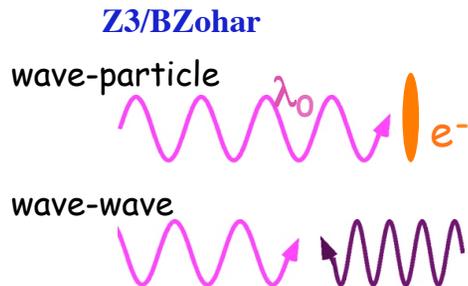


Wave propagation simulations

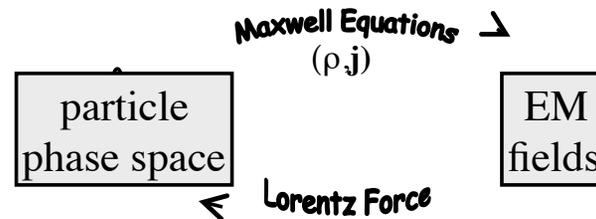


Calculate Reflectivity, Energy Deposition

microscale



Particle in cell (PIC) simulations



Study Saturation Mechanisms

At ~ 10's of petaflops, we can just begin to couple macroscale to mesoscale

Cross-beam energy transfer is in a non-linear, saturated regime (2009 = linear regime)

P. A. Michel

Early NIC results (2009): transfer and symmetry are indeed tunable by $\Delta\lambda$

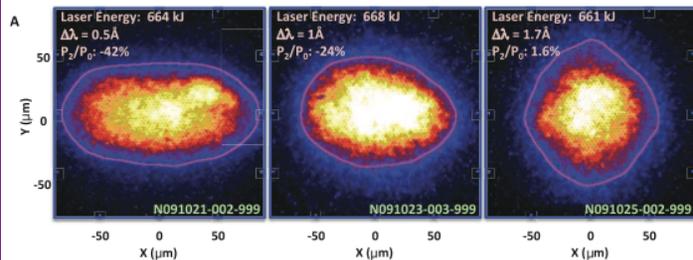
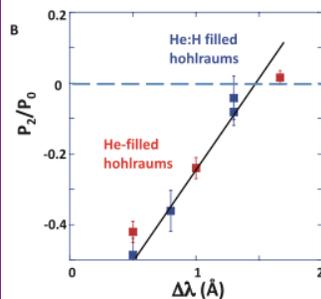


Fig. 3. (A) Capsule x-ray emission images at 9-keV energy from helium-filled hohlraums and 16-ns-long pulses are shown as a function of the wavelength difference between the laser beams on the inner and outer cones. Data are shown at peak x-ray emission at $t = 18 \pm 0.15$ ns. The hohlraum axis of symmetry is vertical. (B) The low-order, even Legendre polynomial is shown at peak emission for helium-filled hohlraums heated with 16-ns-long laser pulses (red symbols) and hohlraums filled with a mixture of helium and hydrogen heated with a shorter, 11-ns-long pulse (blue symbols). The data indicate a linear scaling in agreement with calculations. The error bar is estimated from determining the Legendre polynomial from multiple simultaneously measured images. The dashed curve indicates a symmetric implosion, and the solid curve is the scaling calculated from two-dimensional radiation hydrodynamic simulations.

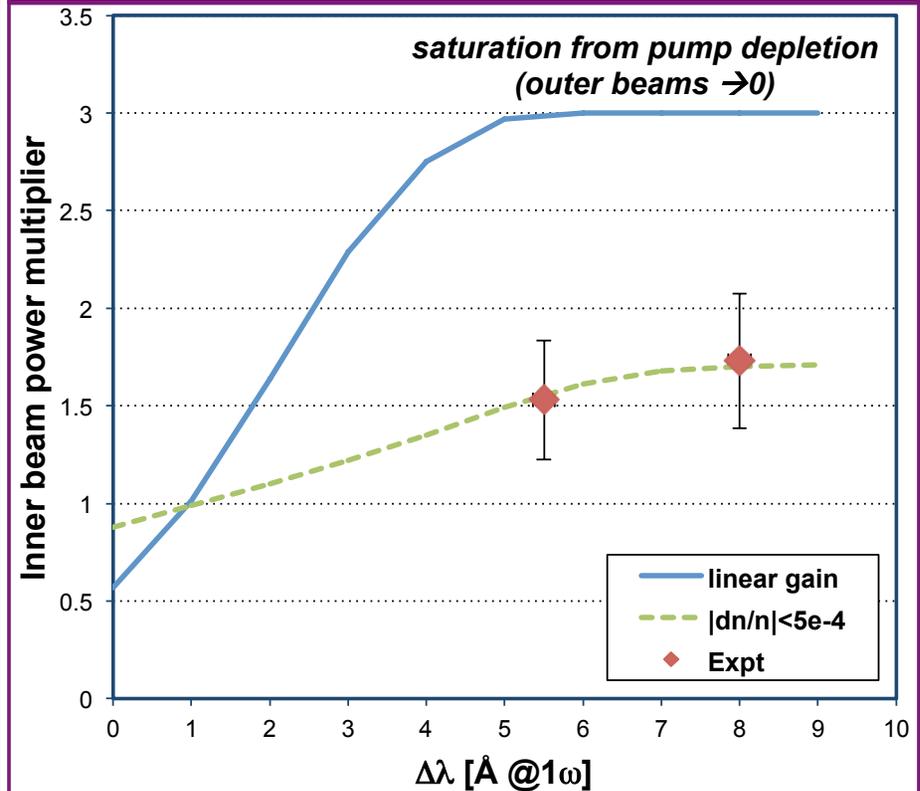


S. Glenzer et al., *Science* (2010)

Symmetry is now exclusively tuned by adjusting cross-beam transfer:

- equatorial symmetry: inner vs. outer beams energy balance ("2-color")
- polar symmetry: 23.5° vs. 30° beams energy balance ("3-color")

2010+: Operation at 1.5 MJ+ and larger $\Delta\lambda$; xbt still robust/tunable, but simulations require non-physical plasma wave saturation

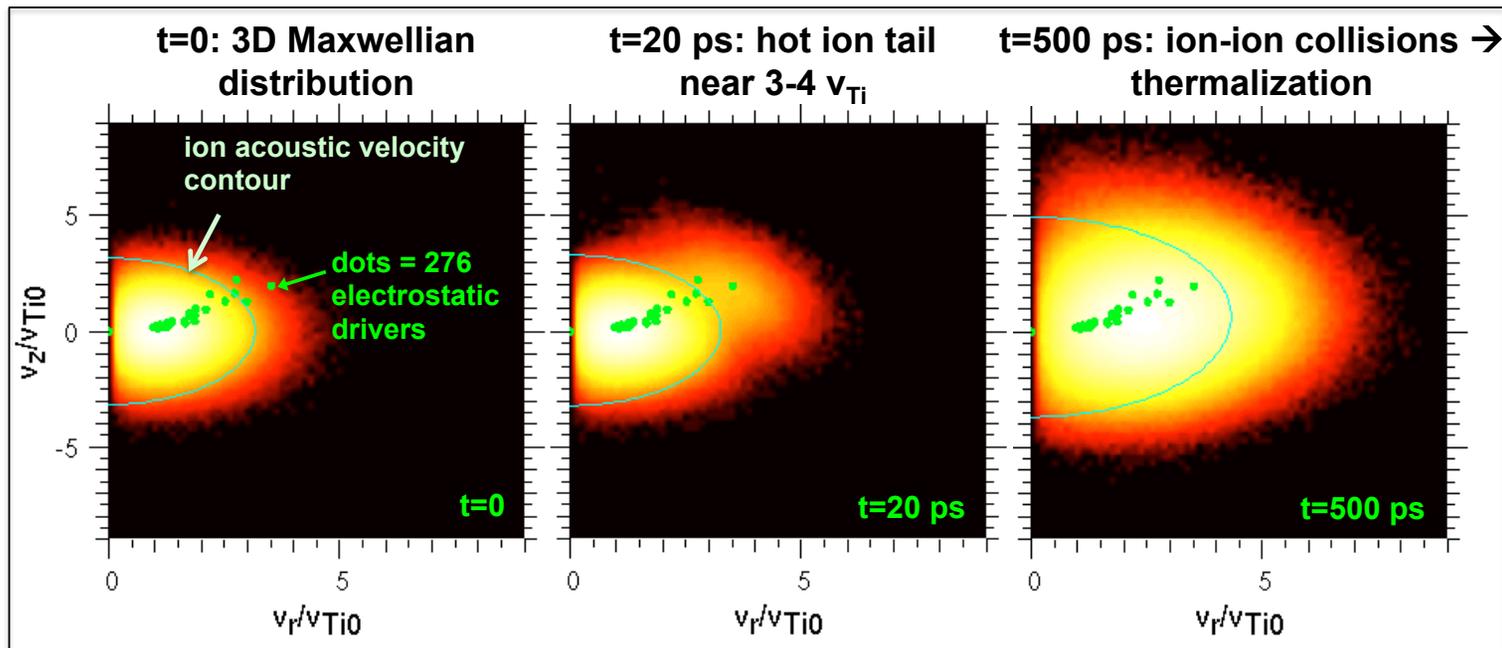


P. A. Michel finds ion heating to be a key saturation mechanism

Calculations show strong stochastic ion heating at the LEH, saturating x-beam transfer (P. A. Michel)

P. A. Michel

At each LEH: 276 possible pairs of quads (i.e. 276 beat waves, i.e. 276 electrostatic drivers): \rightarrow ion acoustic turbulence

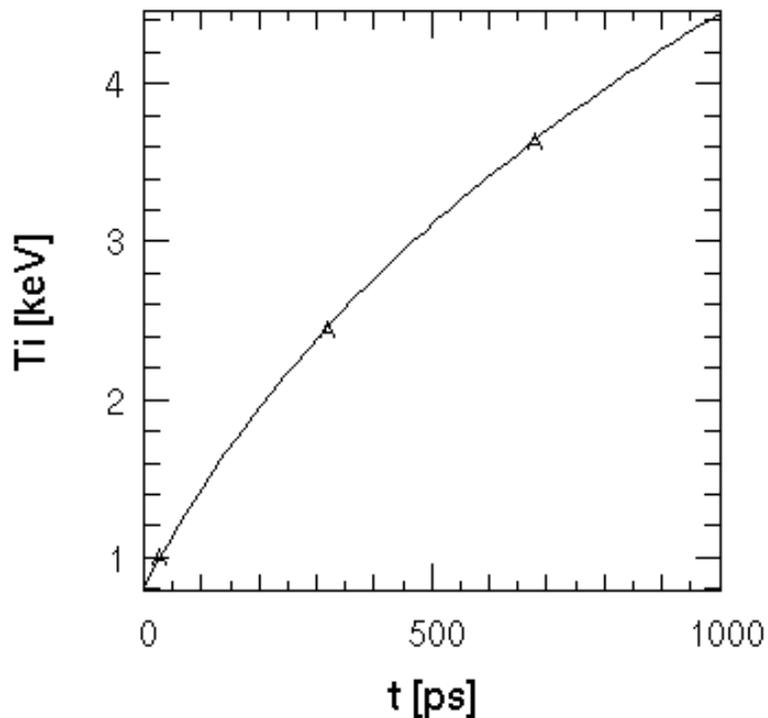


3D “gridless” particle code: self-consistent plasma response given by ion susceptibility via $f(\mathbf{v}, t)$

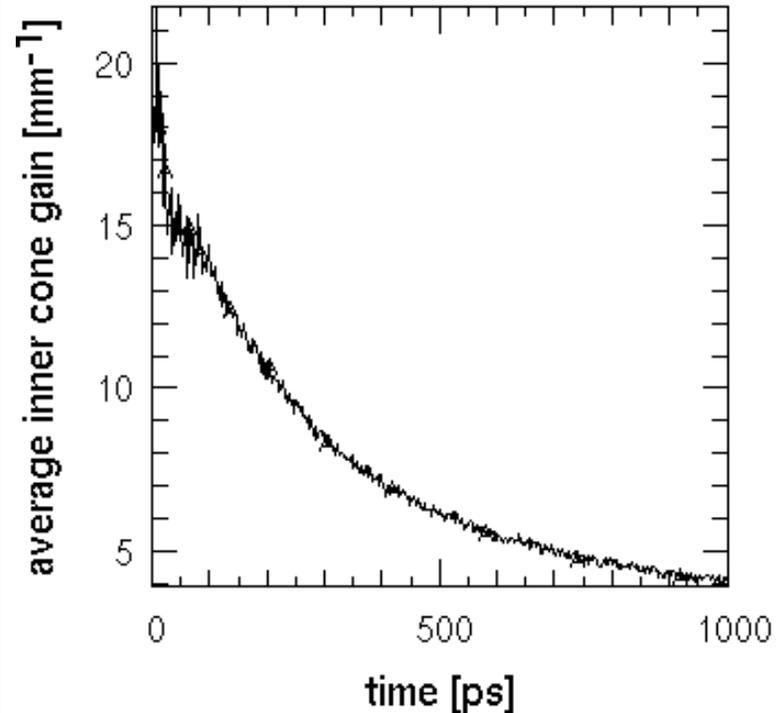
We have developed new numerical tools to calculate the effects of turbulence on cross-beam energy transfer

Stochastic heating \rightarrow strong increase of T_{ion} , which saturates cross-beam transfer over a few 100's ps

Ion temperature increases at a rate of
 \sim few keV/ns



Cross-beam transfer gains decrease by
4-5x over < 1 ns

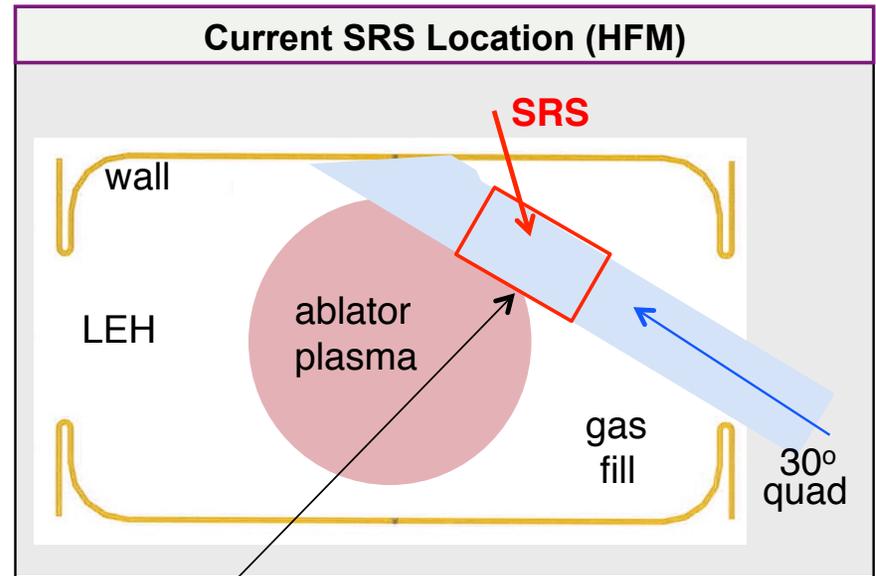
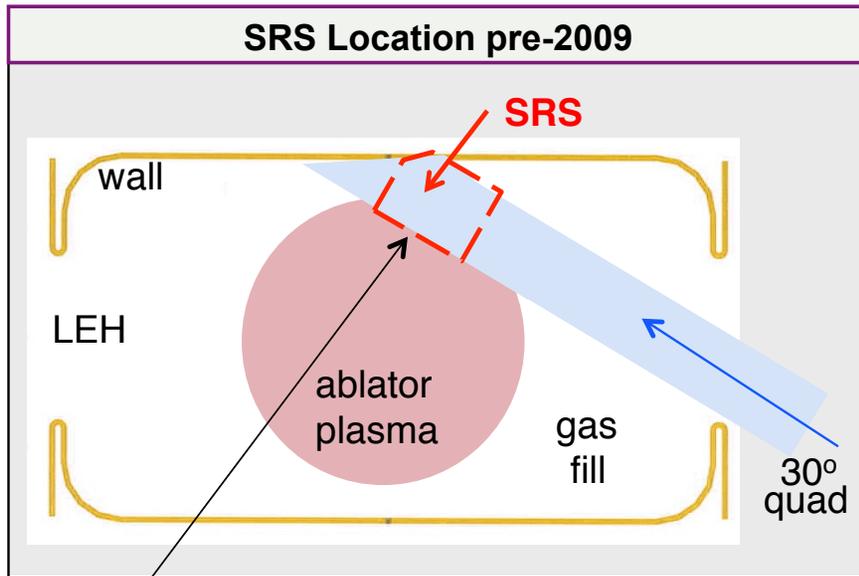


We are working on a self-consistent LPI-hydro implementation of this effect in our hydro codes

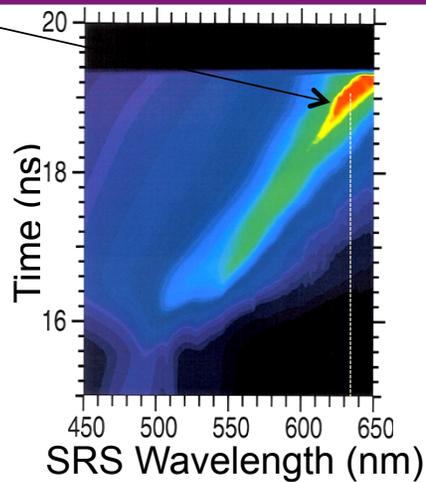
With **laser** and **plasma** modeling improvements, backscatter simulations begin to match experimental data

- Modified plasma conditions from improved hohlraum modeling (high flux model, or “HFM”)
 - results in a cooler plasma – SRS region moves closer to LEH
 - experimental and synthetic SRS spectra show SRS at similar wavelengths
- Realistic laser beams that include the effects of cross-beam energy transfer (XBT)
 - increased power on inner beams (symmetry)
 - spatial non-uniformity in cross section
- Laser quad overlap: simulate beam propagation (pF3D) for two quads of beams (overlap contributes to the intensity; simulate 3-5 quads of beams?)

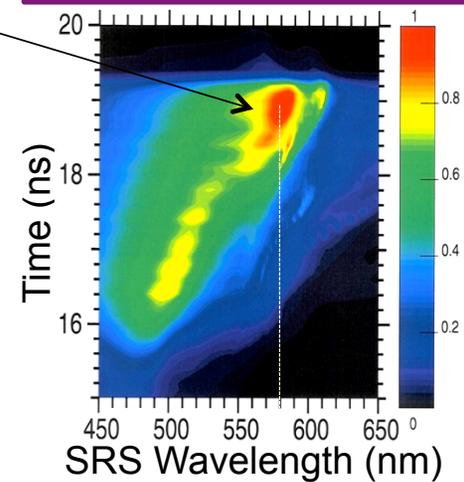
LPI: Hohraum modeling improvements (HFM) have resulted in a change in backscatter location



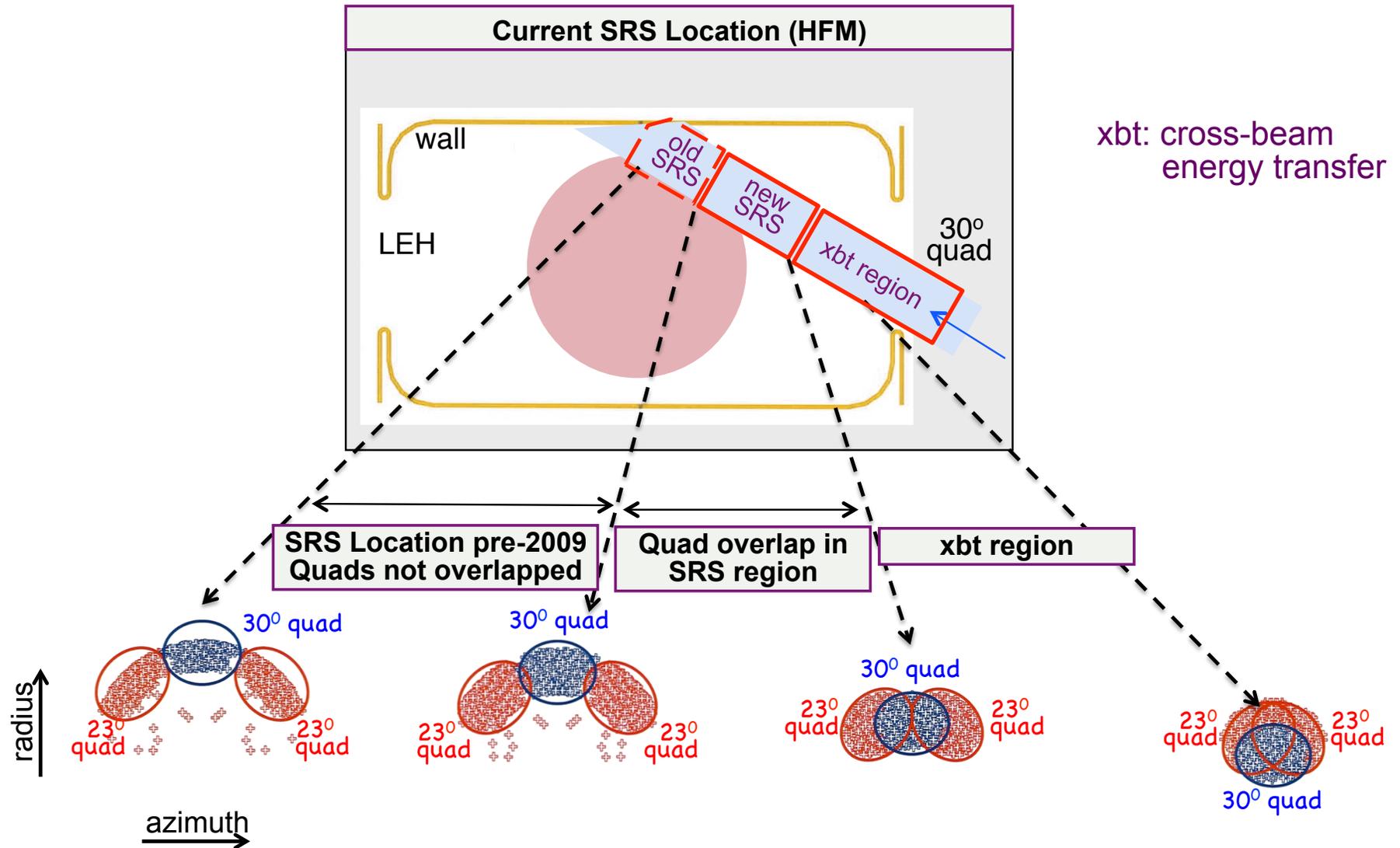
30° SRS Gain Exponent



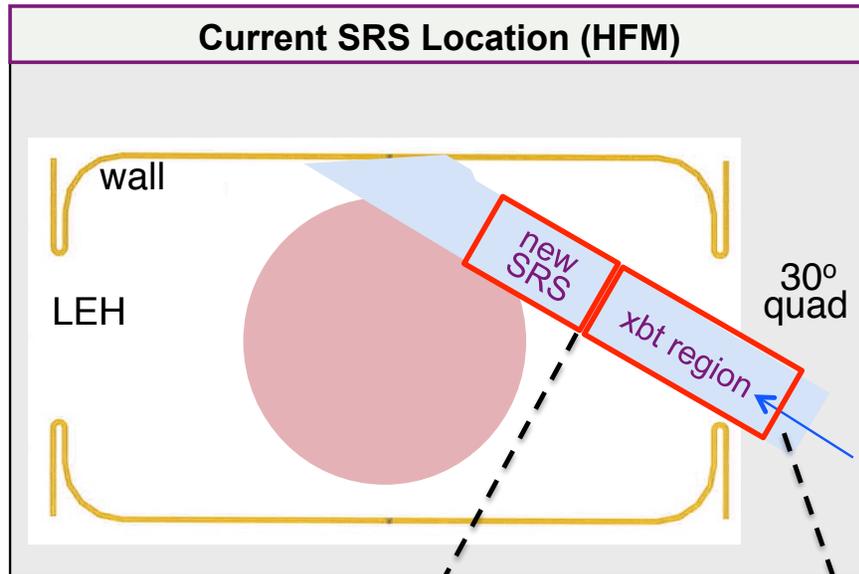
30° SRS Gain Exponent



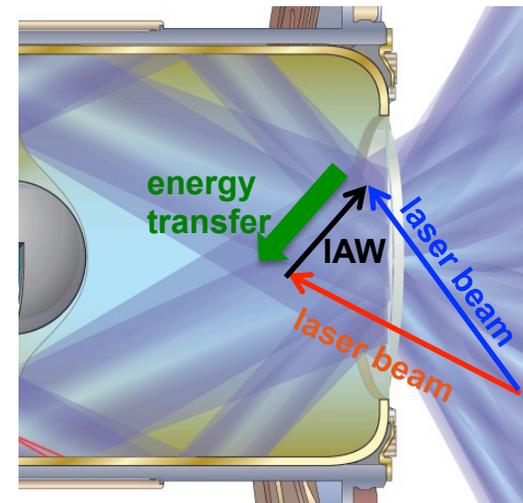
LPI: 30° SRS now occurs where there is overlap with nearest neighbor 23° quads



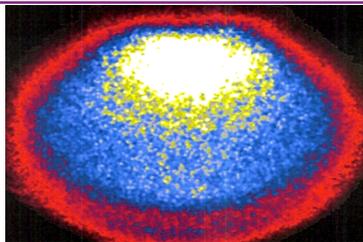
LPI: Cross-beam energy transfer near the LEH results in a spatially non-uniform intensity distribution



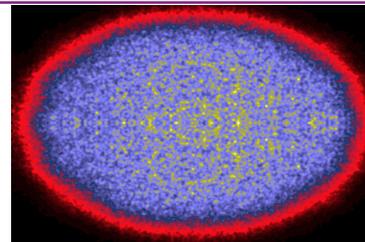
- Cross Beam Energy Transfer (xbt): laser forward scatters off ion acoustic waves (P. A. Michel *et al.*, PoP, May 2010)



Beam after xbt, refraction, absorption



Beam before xbt, refraction, absorption



A promising start: simulated reflectivity approaches experimental levels when these models are used

Experimental Result Vs Single Quad pF3D Simulations

Shot	Energy (MJ)	Time (ns)	30° SRS (TW)
N091204	1.05	19	1.3
2009 pF3D: 1 Quad pre-2009 plasma spatially uniform beam			~ 0.1
2010 pF3D: 1 Quad high flux model plasma spatially uniform beam			0.18
2011 pF3D: 1 Quad high flux model plasma spatially non-uniform beam			0.43

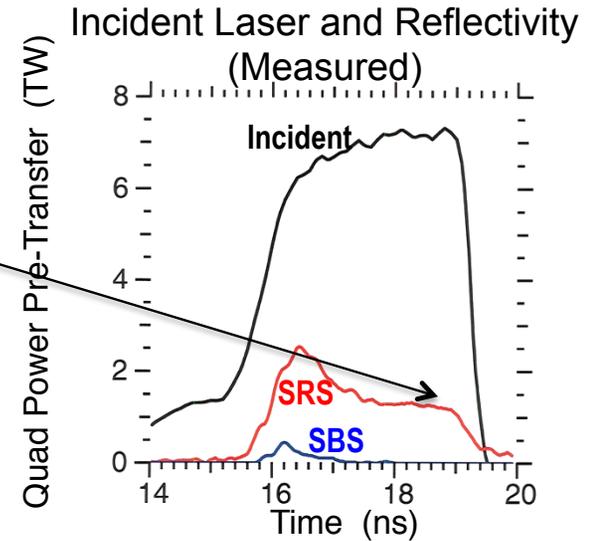
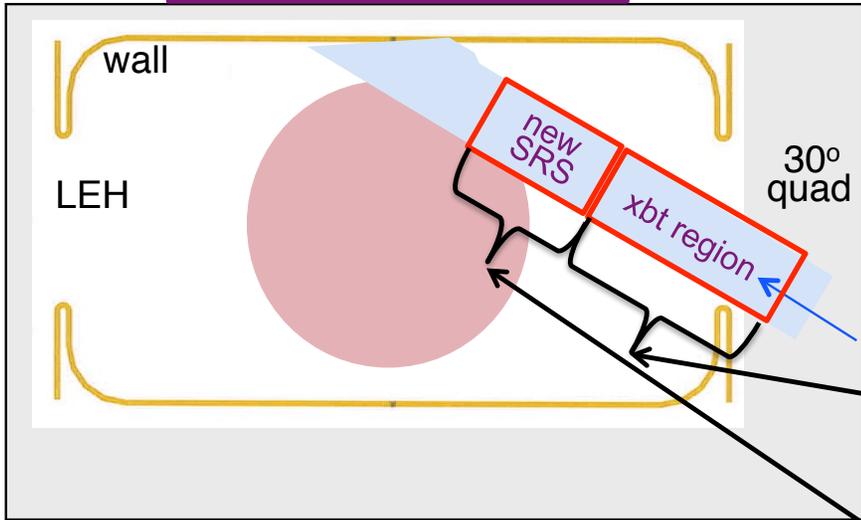
Experimental Result Vs Double Quad pF3D Simulations

Shot	Energy (MJ)	Time (ns)	30° SRS (TW)
N091204	1.05	19	1.3
2010 pF3D: 2 Quads high flux model plasma spatially uniform beams			0.62
2011 pF3D: 2 Quads high flux model plasma spatially non-uniform beams			0.67
2012 pF3D: 3 Quads high flux model plasma spatially non-uniform beams			~0.9-1.0

Simulation in progress on Dawn machine

Proof-of-principle simulation: propagate two quads (23°, 30°) through the resonance region

Simulation @ 19 ns
1.05 MJ Laser Energy
N091204

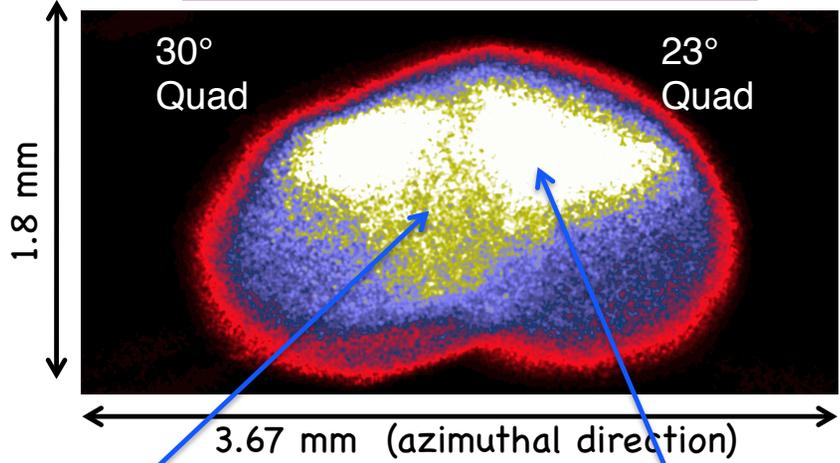


- Laser Input: use SLIP* to propagate quads through the LEH to the input plane (E. A. Williams)
- pF3D: propagate two quads of beams (23°, 30°) through SRS region (A. B. Langdon)

*SLIP code, P. A. Michel and L. Divol

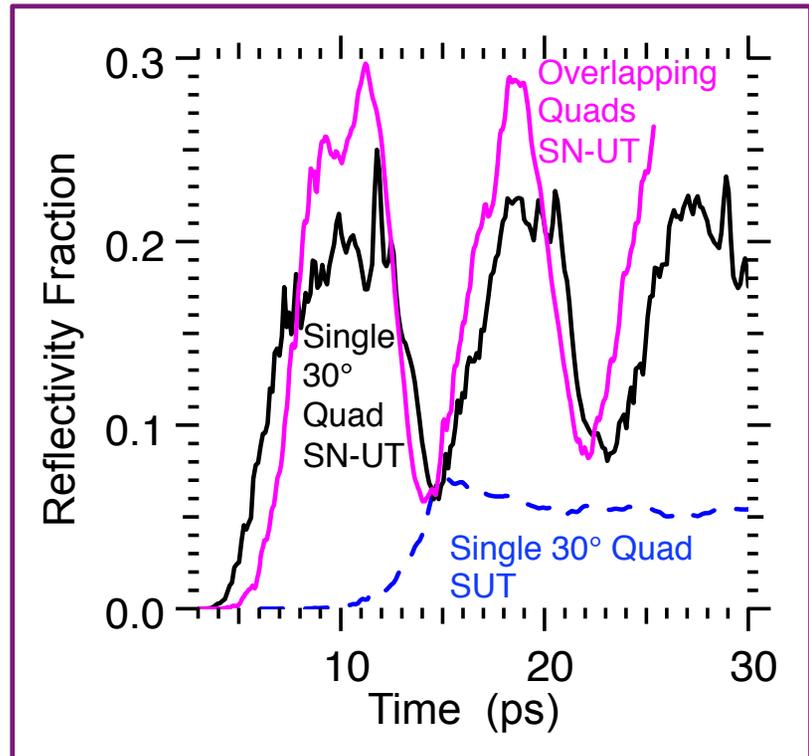
Both quad overlap and spatial non-uniformity increase reflectivity

Laser Quads on Input Plane to "New SRS" Region



overlap region is less intense than xbt region

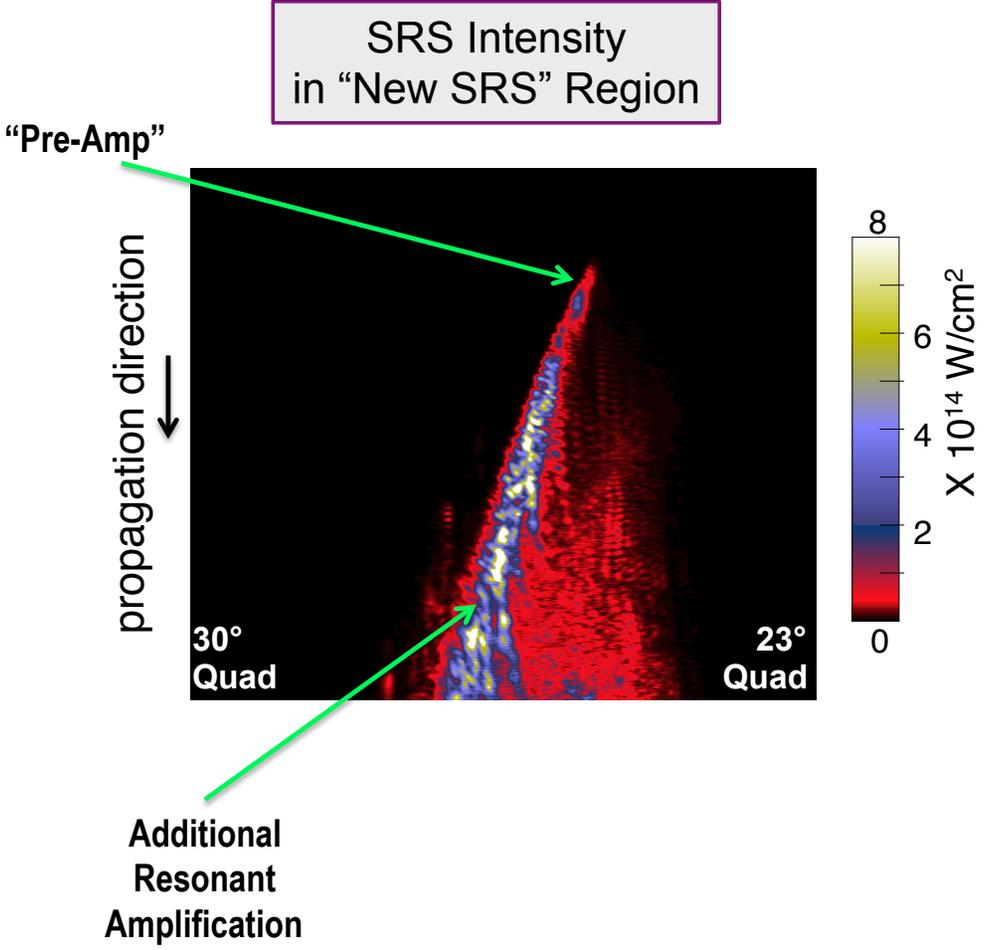
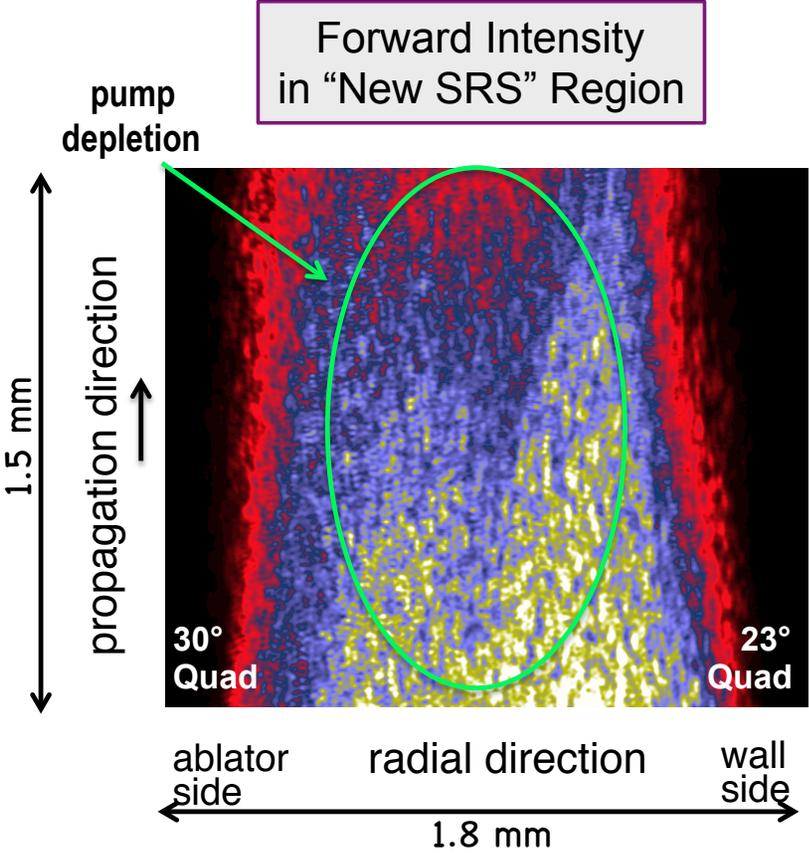
SRS Reflectivity on Input Plane of "New SRS" Region



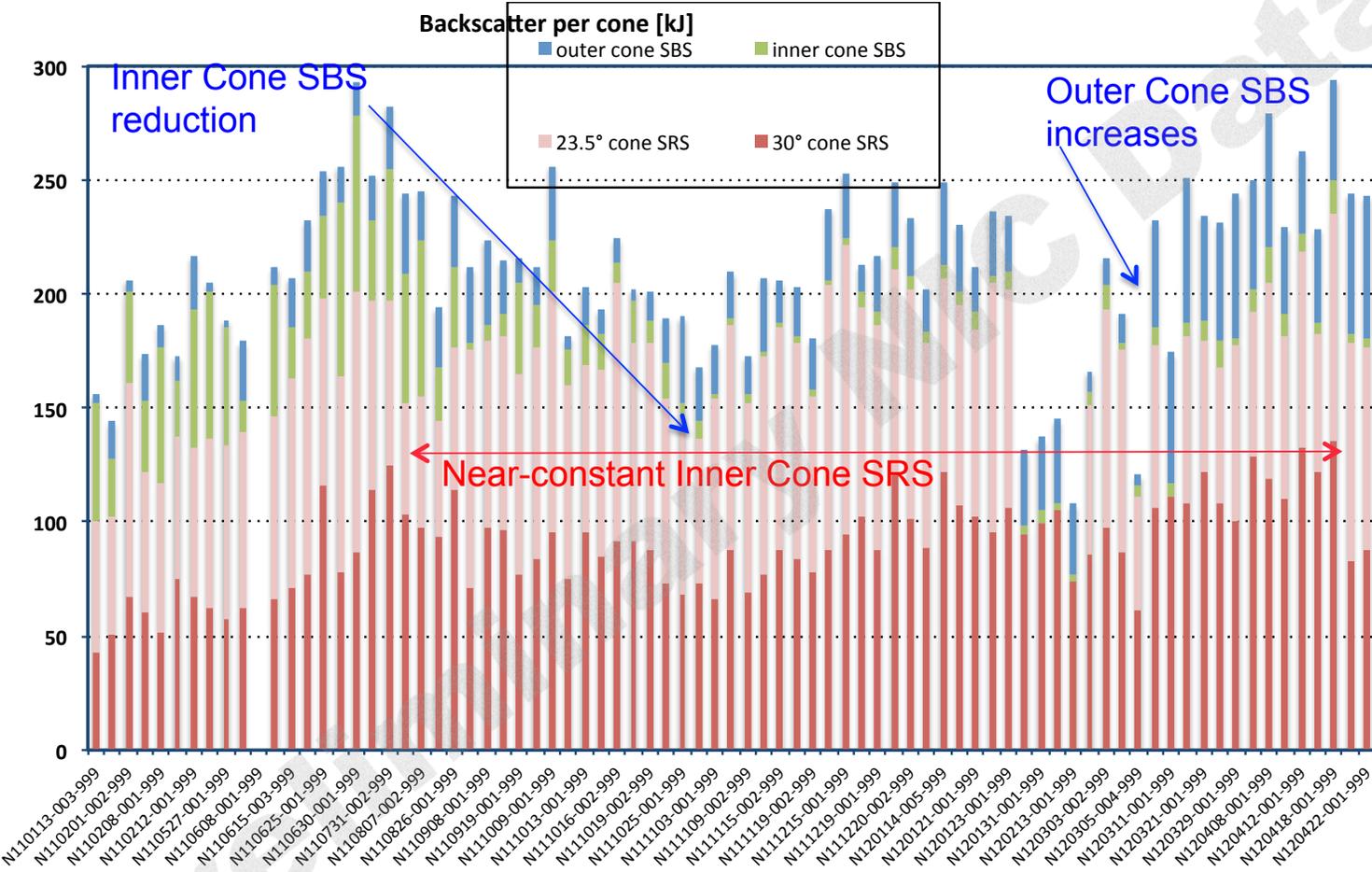
SUT = Spatially Uniform Transfer
 SN-UT = Spatially Non-Uniform Transfer

*We thank the Office of Science INCITE program for computational resources and the ALCF staff for computational support

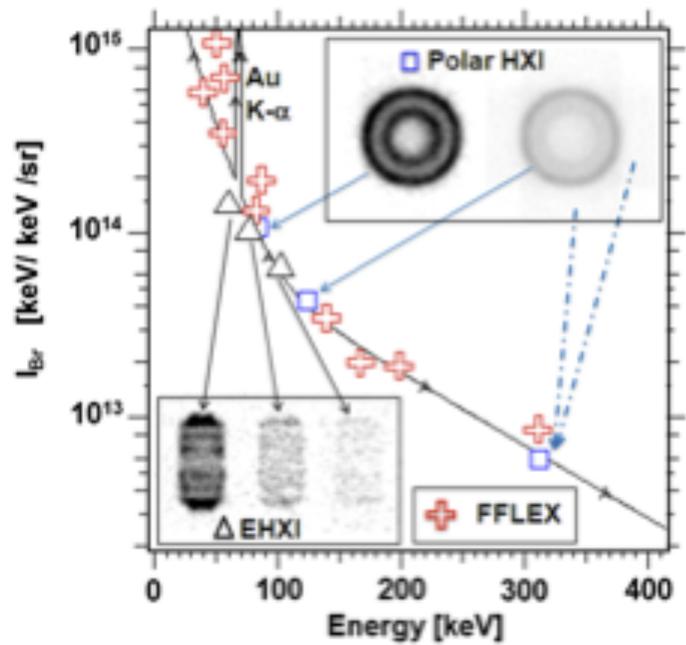
SRS generated by single quads is resonantly amplified in the overlap region



2011-2012 backscatter levels:



A typical NIC hot electron spectrum:





What are the leverage points for improved coupling, and how do we best manipulate them?

$$(1) \quad G = \int_{z_0}^{z_f} dz' \left\{ \frac{k_0^2 v_{osc}^2 \omega_{p2}^2}{v_1 \omega_0 \omega_2} \frac{\Gamma_2}{\Gamma_2^2 + v_2^2 K'^2 z'^2} - \frac{\Gamma_1}{v_1} \right\}$$

0 = incident; 1 = reflected; 2 = plasma wave

$\Gamma \equiv$ damping rate $\omega_{p2} \equiv$ plasma freq.

$v =$ group velocity

$$v_{osc} \equiv c \sqrt{I \lambda^2 / 1.37 \times 10^{18}}$$

CAPABILITIES FOR IMPROVED UNDERSTANDING:

- LPI depends on $n_e, T_e, T_i (\Gamma, \omega_2, \omega_{p2}, K', k_0)$
 - currently inferred from HFM rad-hydro simulations
 - Thomson scatter/spectroscopy provides experimental opinion
 - (i) optical Thomson scatter at LEH;
 - (ii) dope various regions of target?
- Reflectivity in pF3D/VPIC simulations show pumps depletion
 - corroborate with time-resolved (~ 1-2 ps) backscatter diagnostic
- Integration of multi-scale modeling:
 - requires exaflops
- LPI Computational Physicists

(1): E. A. Williams, generalized Rosenbluth gain; applies to backscatter, xbt

What are the leverage points for improved coupling, and how do we best manipulate them?

$$G = \int_{z_0}^{z_f} dz' \left\{ \frac{k_0^2 v_{osc}^2 \omega_{p2}^2}{v_1 \omega_0 \omega_2} \frac{\Gamma_2}{\Gamma_2^2 + v_2^2 K'^2 z'^2} - \frac{\Gamma_1}{v_1} \right\}$$

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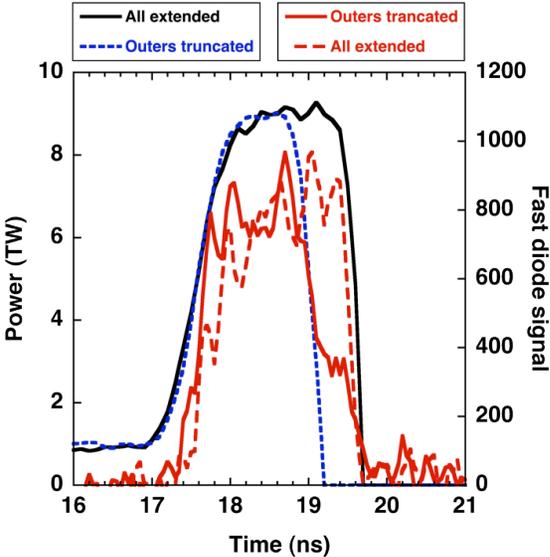
$\Gamma \equiv$ damping rate $\omega_{p2} \equiv$ plasma freq.

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$$v_{osc} \equiv c \sqrt{I \lambda^2 / 1.37 \times 10^{18}}$$

DESIGNS/EXPTS FOR IMPROVED UNDERSTANDING AND COUPLING:

- Design a target to mock up xbt at LEH:
 - measure beam transmission through xbt region (LPI depends on laser intensity)
 - provides information on spatial non-uniformity of beams after xbt
- Design a target that doesn't require xbt for symmetry
 - doesn't require knowledge of beam cross section to understand backscatter
 - might reduce inner cone SRS/hot electrons
 - increase T_e by doping higher-Z material into foam fill?
 - hybrid foam-fill/vacuum hohlraum?
- Design an "SBS-only" target
 - tests our understanding of hot electrons
 - improved coupling?





What are the leverage points for improved coupling, and how do we best manipulate them?

$$G = \int_{z_0}^{z_f} dz' \left\{ \frac{k_0^2 v_{osc}^2 \omega_{p2}^2}{v_1 \omega_0 \omega_2} \frac{\Gamma_2}{\Gamma_2^2 + v_2^2 K'^2 z'^2} - \frac{\Gamma_1}{v_1} \right\}$$

0 = incident; 1 = reflected; 2 = plasma wave

$\Gamma \equiv$ damping rate $\omega_{p2} \equiv$ plasma freq.

$v =$ group velocity

$$v_{osc} \equiv c \sqrt{I \lambda^2 / 1.37 \times 10^{18}}$$

LPI MITIGATION/CONTROL:

- Laser Technology Improvements (affect laser intensity)
 - increased bandwidth
 - STUD Pulses (B. B. Afeyan)
 - increase spot size?
- Reduce scalelength over which SRS/SBS amplifies:
 - introduce density inhomogenities (“zebra foam fills”)
 - introduce temperature inhomogenities (embedded high-Z materials in hohlraum fill?)
- Reduce n_e in hohlraum:
 - shape hohlraum (rugby?) to increase volume between capsule and wall
 - reduces backscatter; less xbt required
- Increase T_e in hohlraum:
 - improves beam propagation
 - reduces backscatter; less xbt required
- Increase damping, reduce n_e , increase T_e :
 - alternate materials for hohlraum fill, ablator, wall

Increased capabilities, dedicated experiments, and testing LPI mitigation will promote improved coupling in ignition hohlraums



<u>CAPABILITIES</u>	<u>EXPERIMENTS</u>	<u>CONTROL/MITIGATION</u>
<ul style="list-style-type: none">• Thomson Scatter (optical, x-ray)• Spectroscopy• Time-resolved backscatter diagnostic• HPC allocation<ul style="list-style-type: none">-- meso/macro-scale cplg-- meso/micro-scale cplg-- expand at mesoscale• Dedicated computational physicist to LPI	<ul style="list-style-type: none">• Transmitted beam thru xbt region• Hohlraum that doesn't require xbt for symmetry• Hohlraum that only generates SBS	<ul style="list-style-type: none">• LASER:<ul style="list-style-type: none">-- bandwidth-- STUD pulses• PLASMA:<ul style="list-style-type: none">-- increase T_e (add high-Z fill dopant)-- reduce density (change hohlraum shape)-- introduce inhomogeneities ("zebra" foam fills)-- improve beam propagation, increase damping (alternate ablators, fills, walls, high-Z dopants)

NIC

