

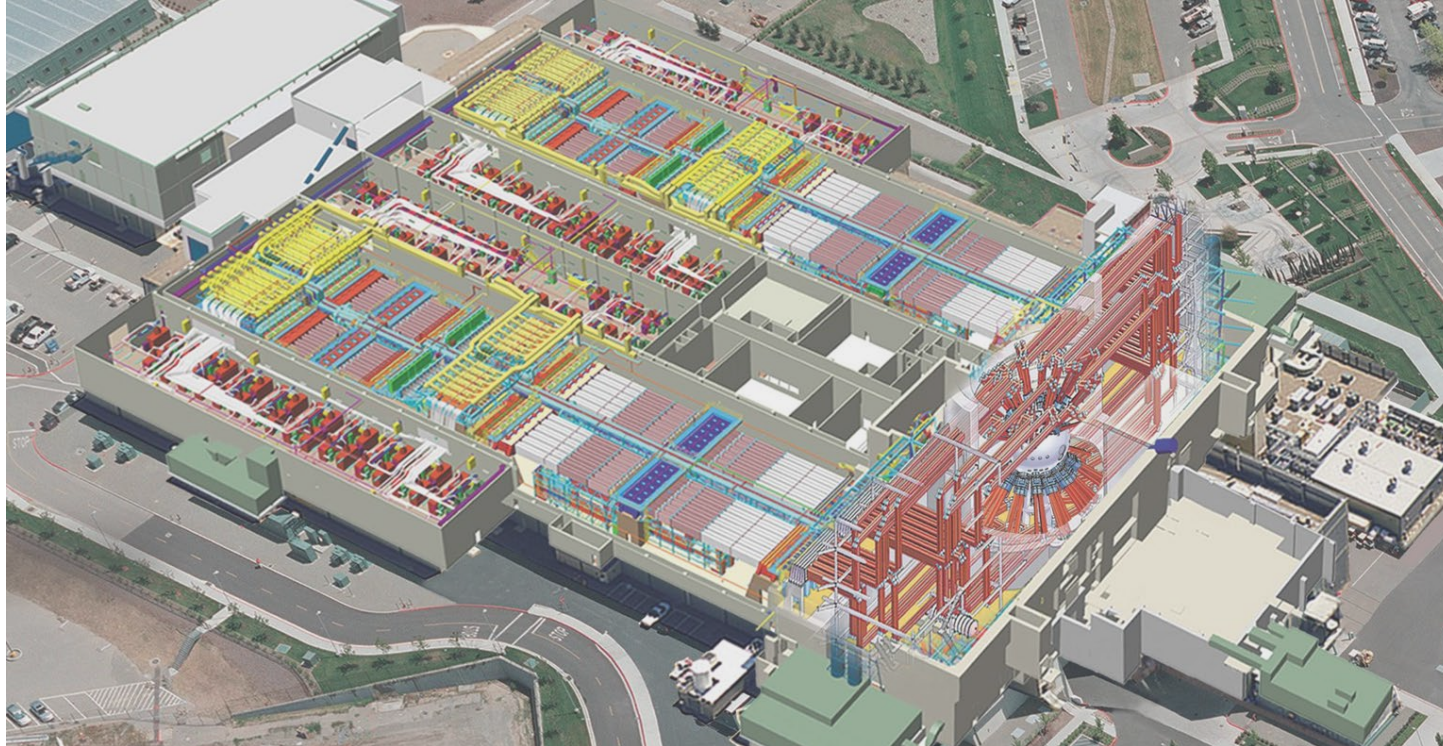
# Present understanding of ignition and Gain using indirect-drive inertial confinement fusion target designs on the U.S. National Ignition Facility

Omar A. Hurricane, Chief Scientist, ICF Indirect-Drive Collaboration

NIF/JLF Users Group  
Garré, February, 11 2025

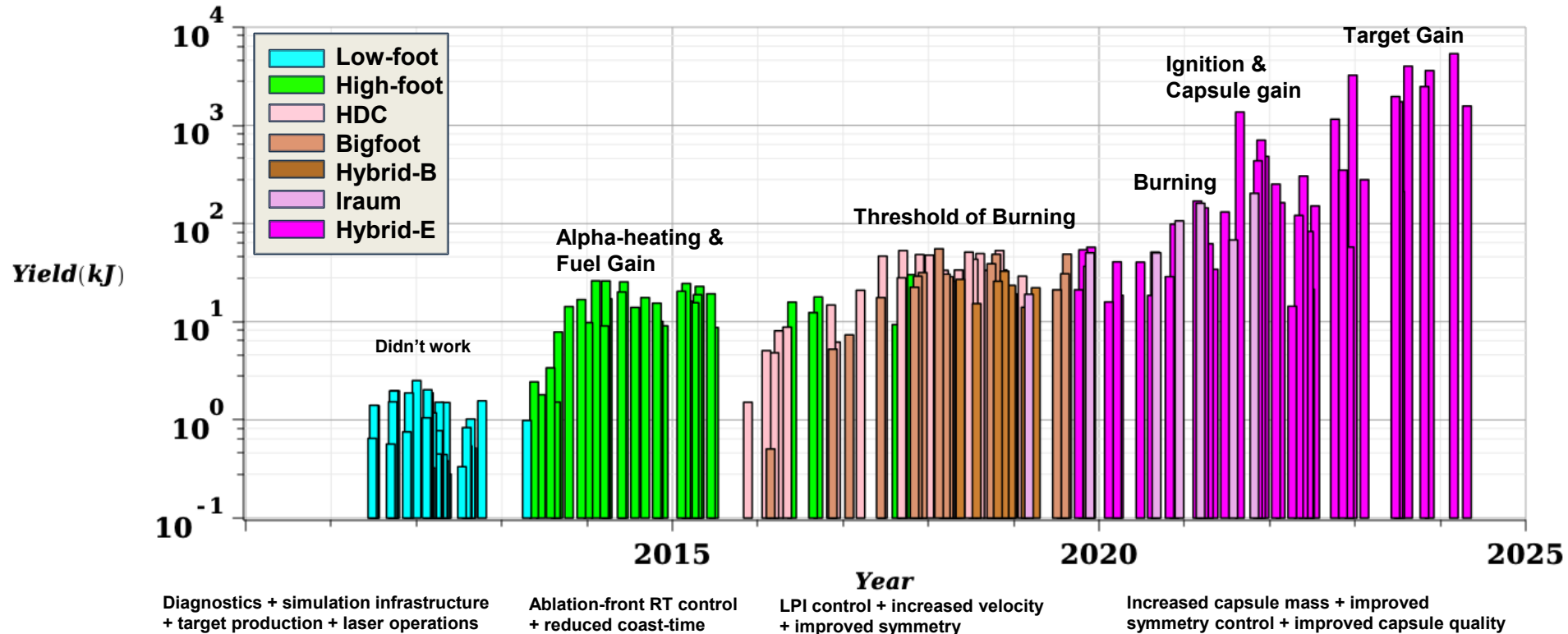


# The National Ignition Facility (NIF) is the first laboratory fusion fusion facility in the World to demonstrate fusion “ignition”

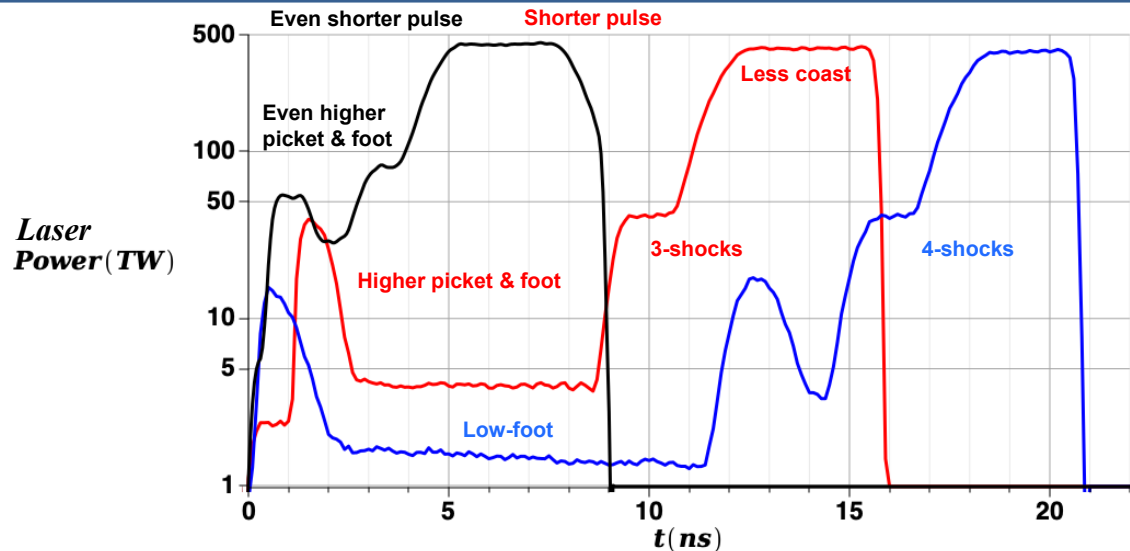
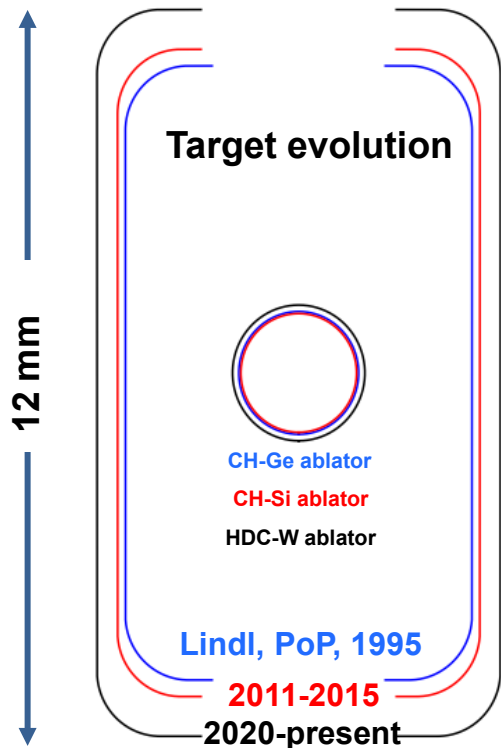


- Size of 3 football fields
- 327 MJ of stored electrical energy
- 2.2 MJ of 351 nm (“blue”) light directed into the target chamber
- 20 years to design and build
- 10 years of experiments to get ignition

# We have an “existence proof” that ignition in the laboratory is possible, but getting ignition has been extremely difficult



# Today's target design\* on the NIF was built upon a 10+ year *evolutionary exploration* of how to manage all the technical challenges

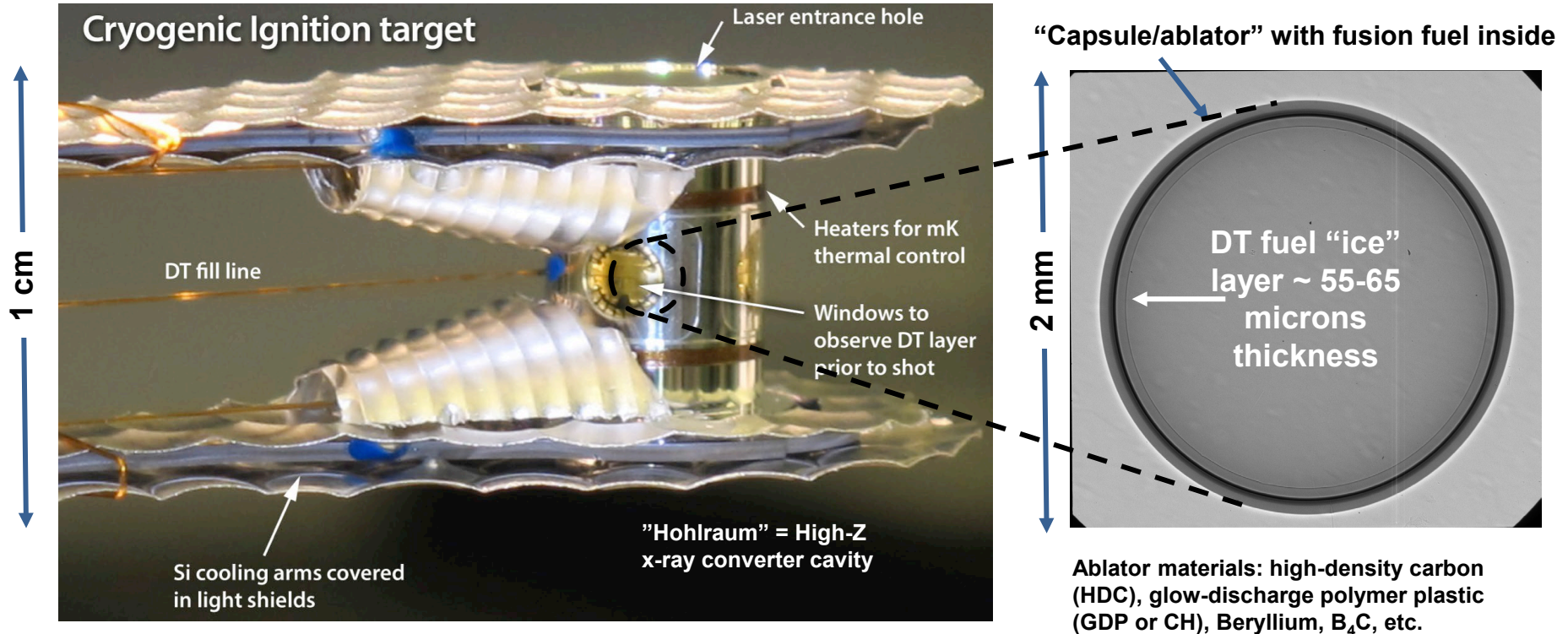


**Qualitative Design: Unchanged**

**Strategy: High energy gain → Just get something to ignite, even if barely**

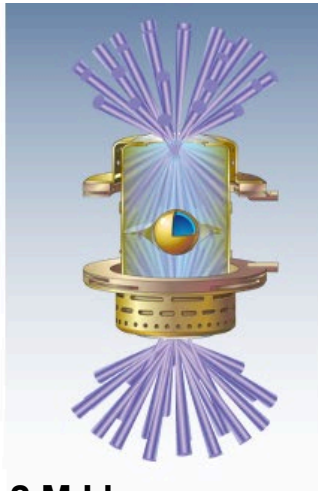
**Tactics: Many detailed changes – extraordinary sensitivity to “small” changes**

# Inertial confinement fusion uses highly engineered “targets” that are complicated, fragile, and (because of reality) imperfect



# Indirect drive inertial confinement fusion (ICF) uses x-rays to ablate and accelerate a capsule of fusion fuel to extreme velocity

Lasers deposit energy into hohlraum



2 MJ laser energy

A bath of x-rays is created as the hohlraum heats

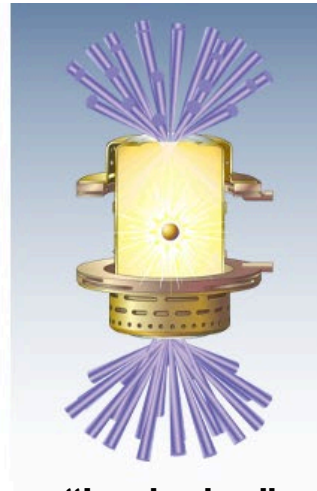


The capsule surface ablates at ~150 Mbar



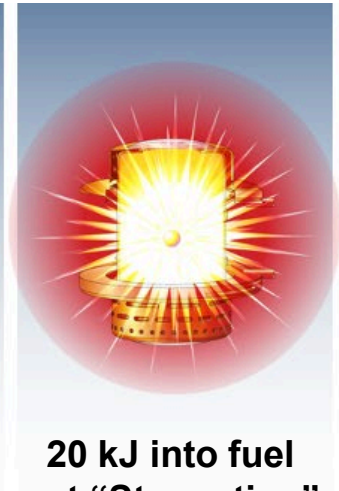
250 kJ absorbed

Capsule accelerates inwards doing  $p dV$  work



“Implosion”

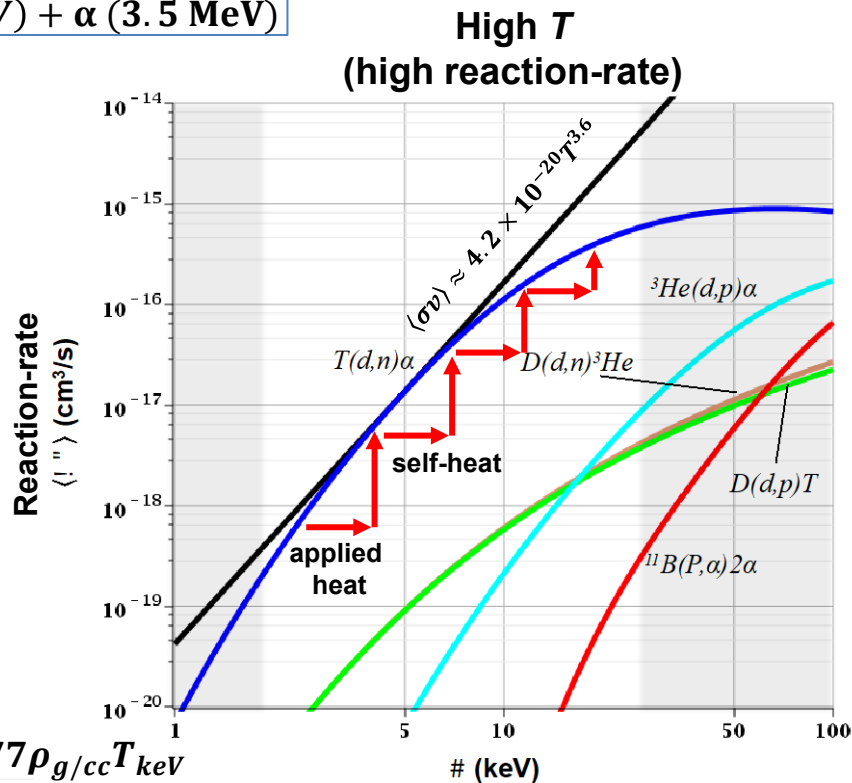
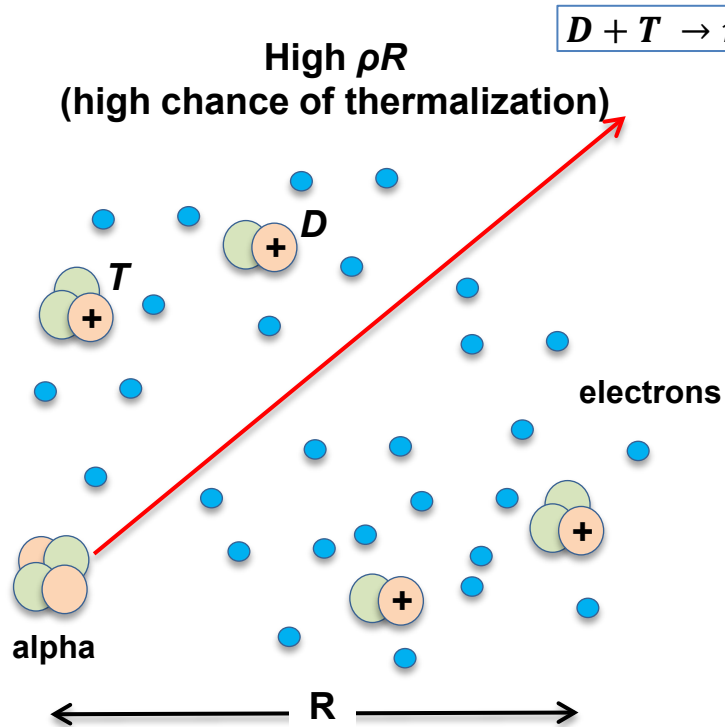
Kinetic energy is converted into internal energy



20 kJ into fuel at “Stagnation”

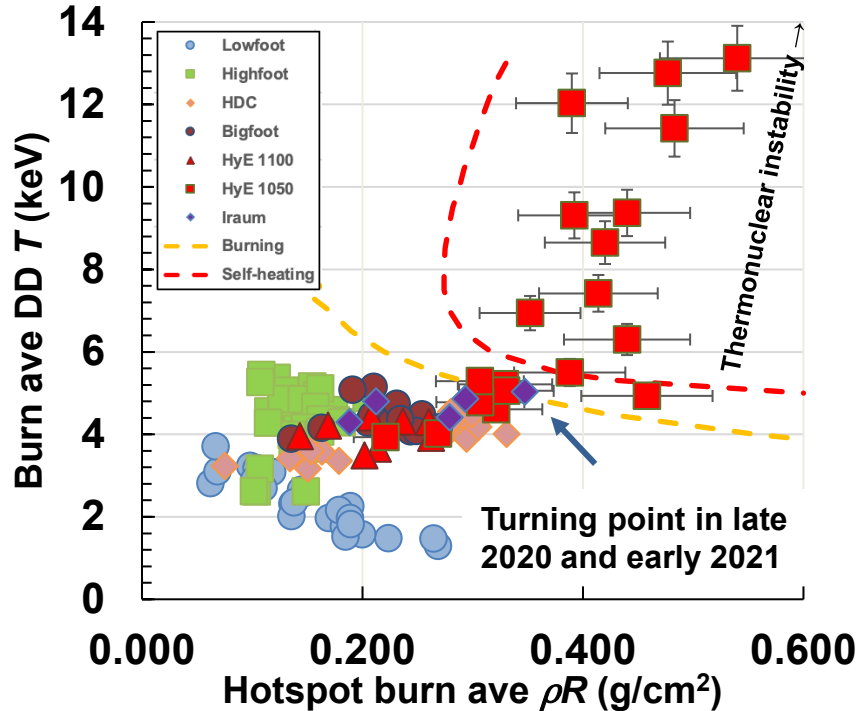
Energy is sacrificed to amplify pressure in the fusion fuel,  $pR^5 \sim \text{constant}$

# As the alpha's deposit energy inside the plasma *thermonuclear instability* ("ignition") can occur if $\rho RT$ is large enough for long enough



$$P_{DT}(\text{Gbar}) \sim 0.77 \rho_{g/cc} T_{\text{keV}}$$

# Obtaining the thermodynamic conditions for fusion ignition was exceeding difficult, but the theoretical threshold was where we expected it to be



Example  $\alpha$ -stopping:

$T \sim 10 \text{ keV}$  ( $\sim 110,000,000 \text{ Kelvin}$ )  
 $\rho R \sim 0.4 \text{ g/cm}^2$

$\sim 70\%$  of alpha's deposit energy  
 in hot plasma

"Once a burning plasma was achieved, ignition wasn't a question of *if* but *when*." R. Betti, *Nature* 2022

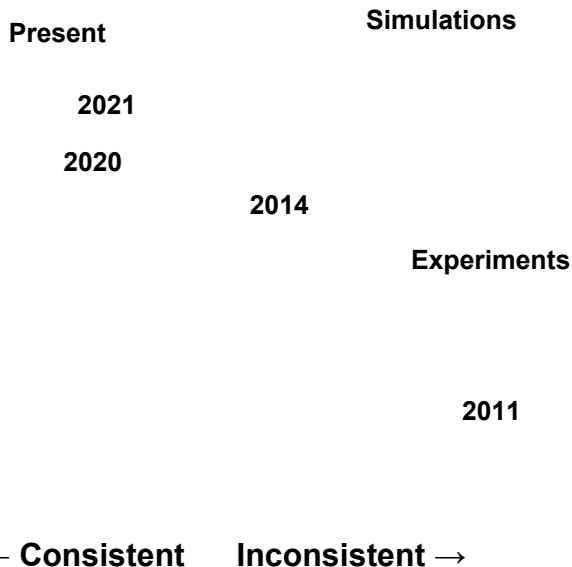


# Early ignition experiments on the NIF tried for ultra-high compression, but we had to back-off to get things working

## NIF DT experiments fuel gain vs. fuel areal density

Fusion yield/energy into the DT

From Hurricane, et al., Plasma Physics and Controlled Fusion. Phys, 2024



**Change of strategy:**

**Instead of going straight for high yield and high gain, back-off and find a regime where simulations and data start to better agree.**

**Then move forward by solving problems in steps.**

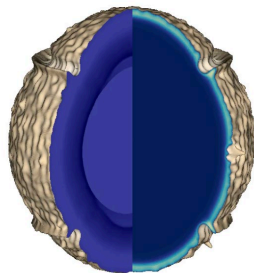
**“Basecamp strategy”**

**We still don't understand all the reasons for the discrepancy**

# Implosions can be nonlinear *problem-amplifiers*, amplifying the hydrodynamic perturbations of defects and asymmetry

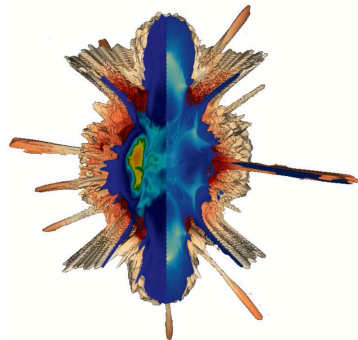
## Instability & defects:

e.g. D. Clark *et al.*, Phys. Plasmas 23, 056302 (2016)

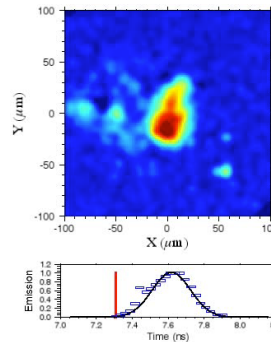


$T_{ion}$        $\rho$

Time →



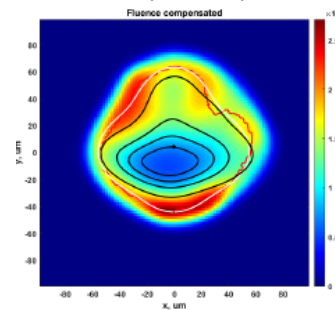
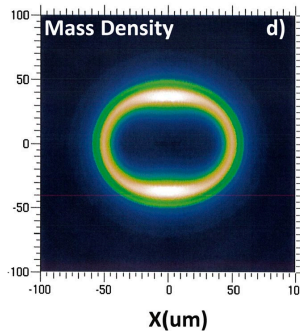
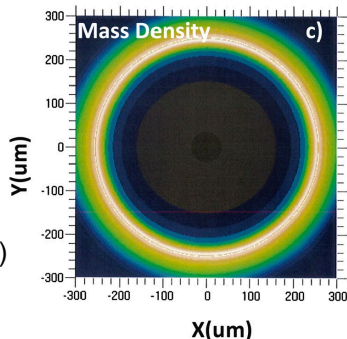
$T_{ion}$        $\rho$



FC Downscatter Data (N210307):

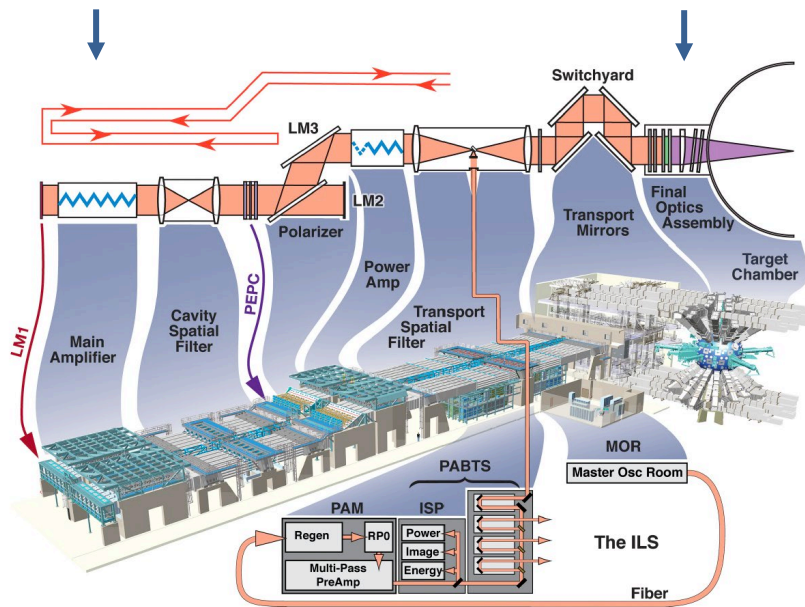
## Asymmetry:

e.g. A. Kritcher *et al.*, Phys. Plasmas 21, 042708 (2014)



Both problems effectively "cost" us energy

# In the future, the NIF will extend\* the present 2.2 MJ laser energy capability, at 351 nm, to 2.6+ MJ for use with larger targets



M. L. Spaeth, et al, Fusion Science & Technology, 69, 25-145, 2016

\*J. M. Di Nicola, Nuclear Fusion, 59, 032004, 2019

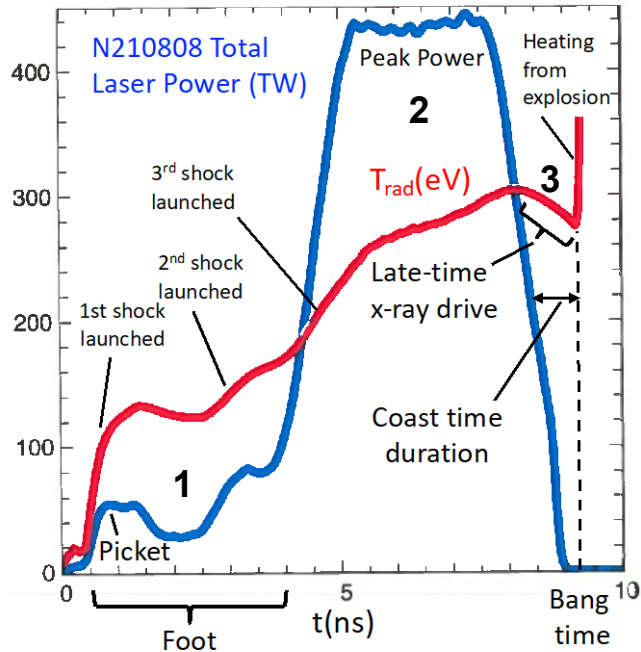
# Remarkable that we can now talk about burning plasmas, ignition, and scientific breakeven ( $G_{\text{target}} > 1$ ) in the past-tense!

- **Maximizing energy delivery to the DT fuel without pressure loss has been key**
  - Asymmetry costs *energy*
  - Instability and mixing costs *energy*
  - Long coast-time costs *power*
- **“Ignition” is where we expected it in terms of plasma conditions, but engineering control (of laser and targets) was/is extremely challenging and more laser energy than expected was needed**
- **While “scientific breakeven” was obtained on December 5, 2022, the NIF laser facility consumed  $\sim 100 \times$  more energy (327 MJ) than was created with DT fusion**
- **Future work involves extending the laser energy delivery capability of the NIF to drive larger scale targets to higher yields and to try and solve some of the outstanding mysteries we’ve encountered over the past 12 years of work**

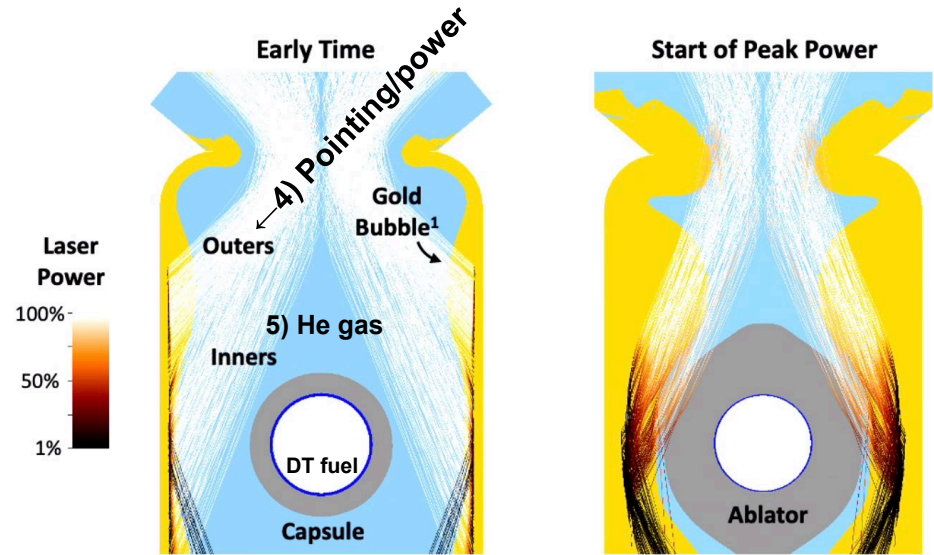
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# The NIF laser in combination with an x-ray driven target creates an implosion that compresses and heats the fusion fuel



- 1) Stability & entropy (“adiabat”)
- 2) Peak velocity
- 3) Coast-time (mechanical power)



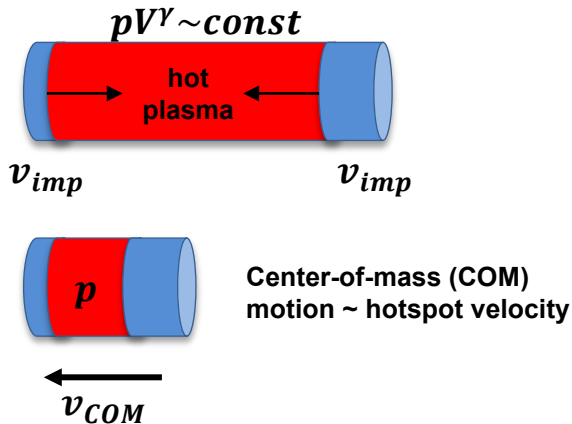
Ablation pressure:  $p_{abl} \sim T_{rad}^{2.5-3.5}$

- 4) Laser pointing and “cone” power (symmetry)
- 5) Helium gas-fill (resists plasma ingress)

# Implosion symmetry control is important, because it wastes shell KE, that could have heated & compressed the fusion fuel

## Asymmetric implosion abstracted to pistons

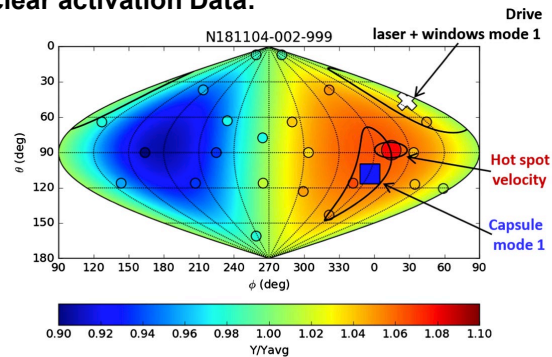
Mode-1  
classical  
mechanics  
model



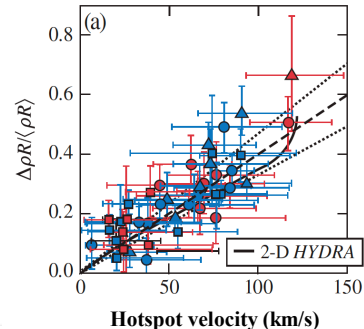
Unrecoverable KE due to COM motion:

$$f^2 \equiv \left( \frac{\rho \delta R_{max} - \rho \delta R_{min}}{\rho \delta R_{max} + \rho \delta R_{min}} \right)^2 = \frac{v_{com}^2}{v_{imp}^2} = nRKE$$

## Nuclear activation Data:

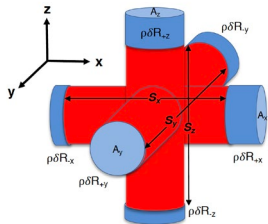


## Nuclear time-of-flight (NTof) Data:



# 3D geometry reflects energy efficiency, and one finds the key measure is the weighted harmonic mean of shell $\rho\delta R$

3D abstraction:



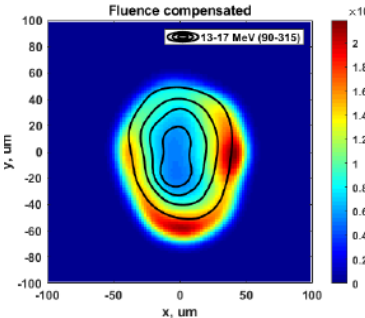
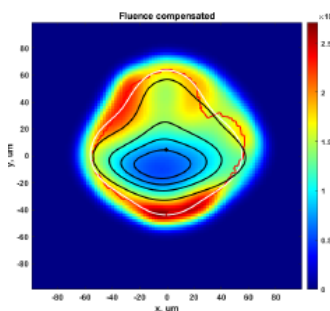
Newton's laws  $\rightarrow$

$$f^2 = 1 - \frac{\left\langle \frac{1}{\rho\delta R} \right\rangle^{-1}}{\langle \rho\delta R \rangle}$$

Weighted harmonic mean of shell areal density

where  $\langle \blacksquare \rangle = \frac{\int \blacksquare dA}{\int dA}$  and the hotspot surface area weighting is  $dA$

FC Downscatter Data (N210307):



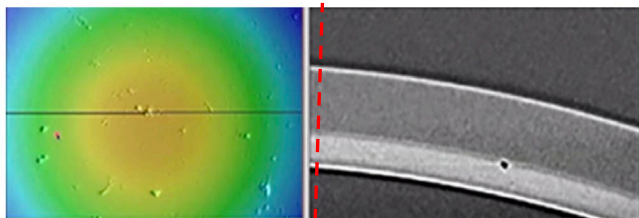
FC Downscatter Data (N201101):

Thin regions reduce implosion efficiency

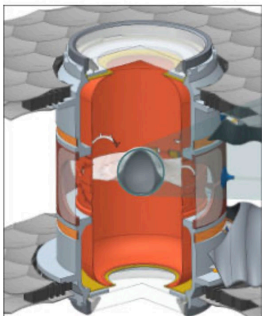


# Defects in targets can have a significant energetics consequence and can be a limiting factor for ignition and energy gain

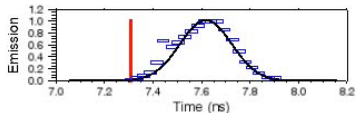
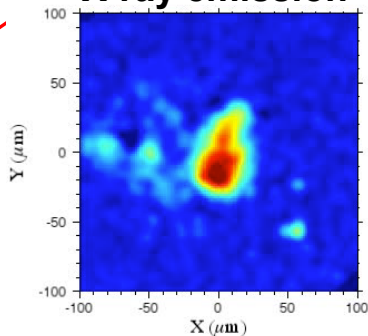
## Capsule defects seed instability and mixing



(Left) Confocal microscope image showing pits on the capsule surface. (Right) Tomographic image showing an internal void.

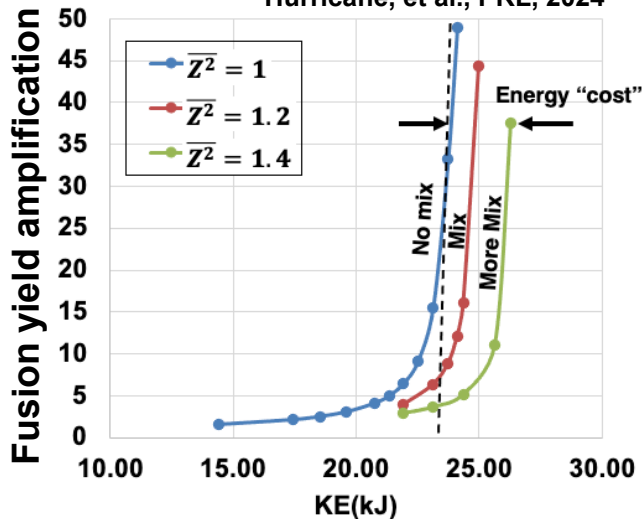


### X-ray emission



Braun, et al., Nuclear Fusion, 63, 2022

Hurricane, et al., PRL, 2024



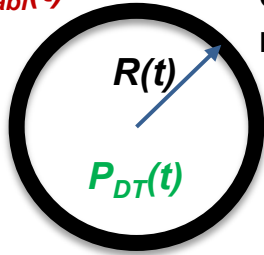
Ignition threshold moves to higher energy:

$$KE_{mix} \approx KE_{no-mix} \overline{Z^2}^{0.3}$$

# ICF implosions are mechanical systems that act as nonlinear power and pressure amplifiers, as kinetic energy converts to internal energy

Ablation pressure

$P_{abl}(t)$

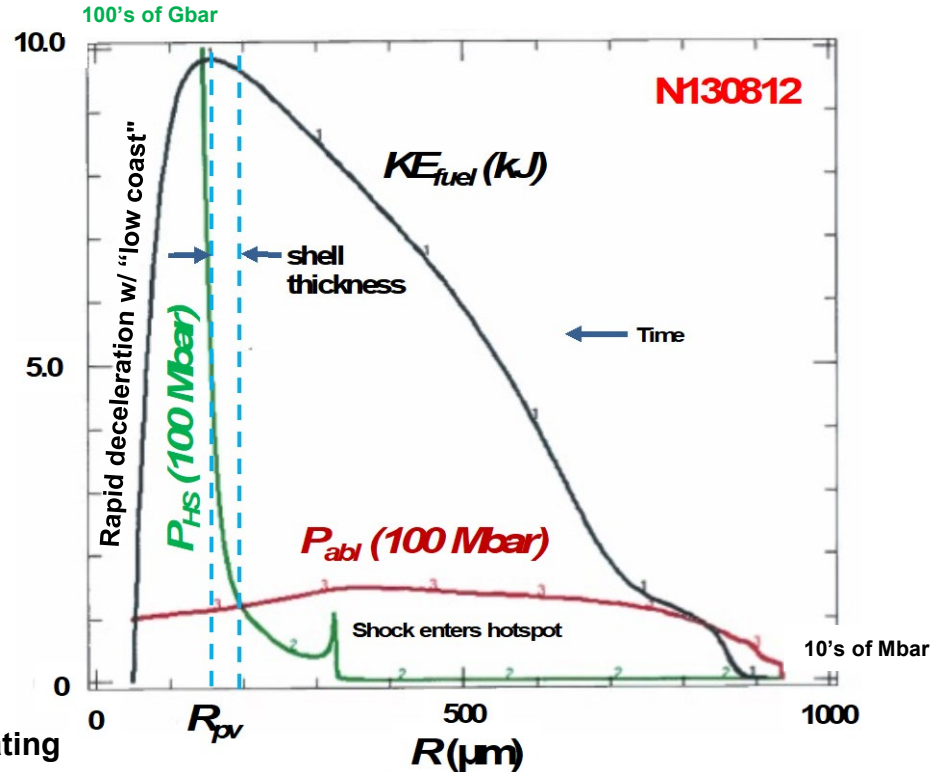


Shell (fuel + ablator)  
mass,  $m_{sh}$

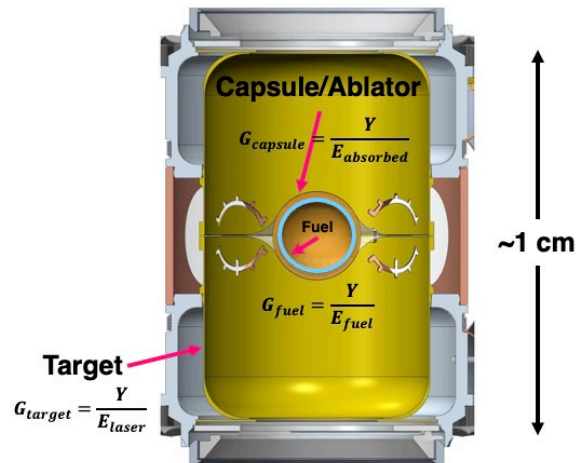
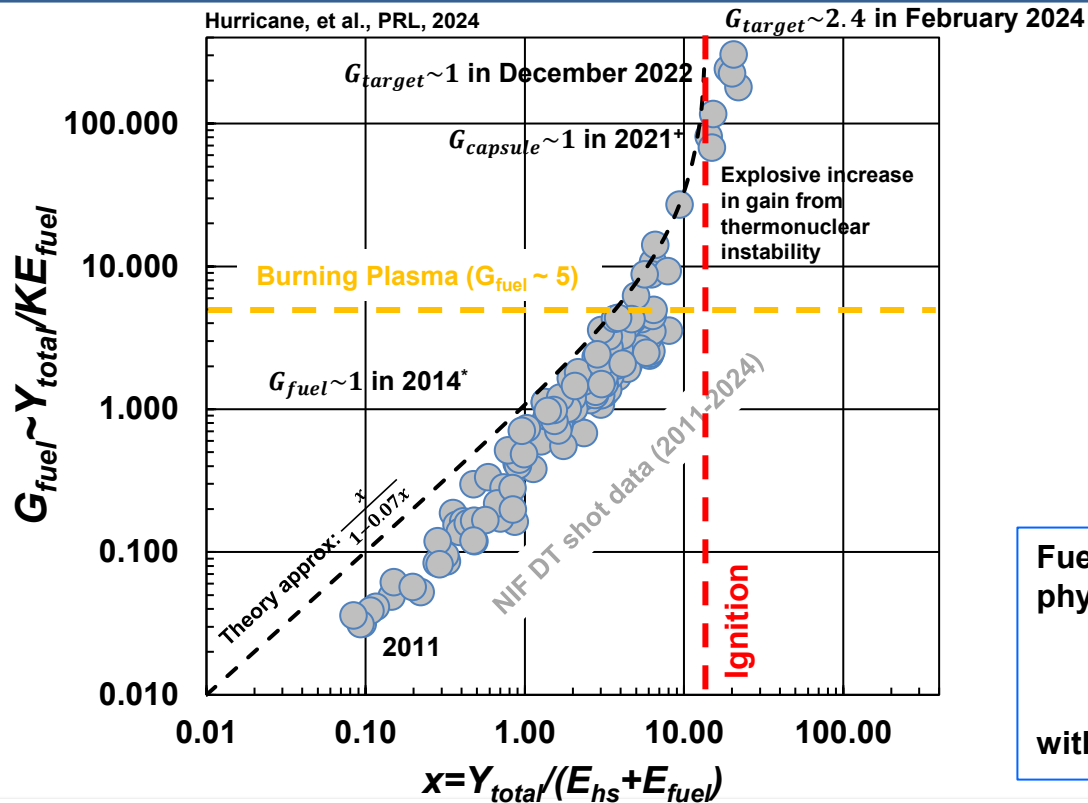
Thermodynamics:  $P_{DT} \approx P_{pv} \left( \frac{R_{pv}}{R} \right)^5 e^{P_{DT} \tau H(T, \bar{Z})}$

amplification from convergence

implicit exponential  
amplification from alpha-heating



# Three energy “gain” (energy out/energy in) milestones have been achieved over that same decade of work

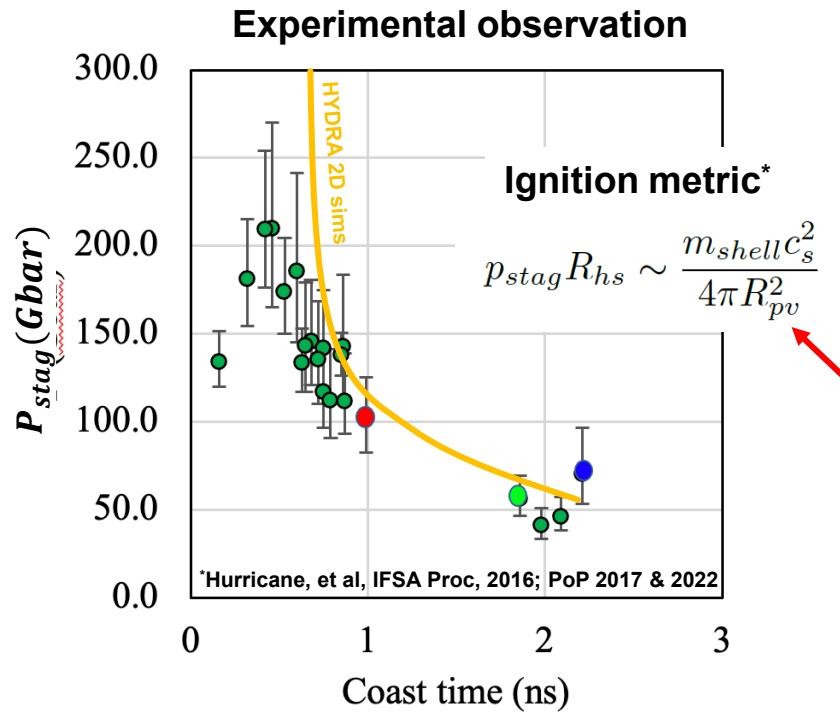
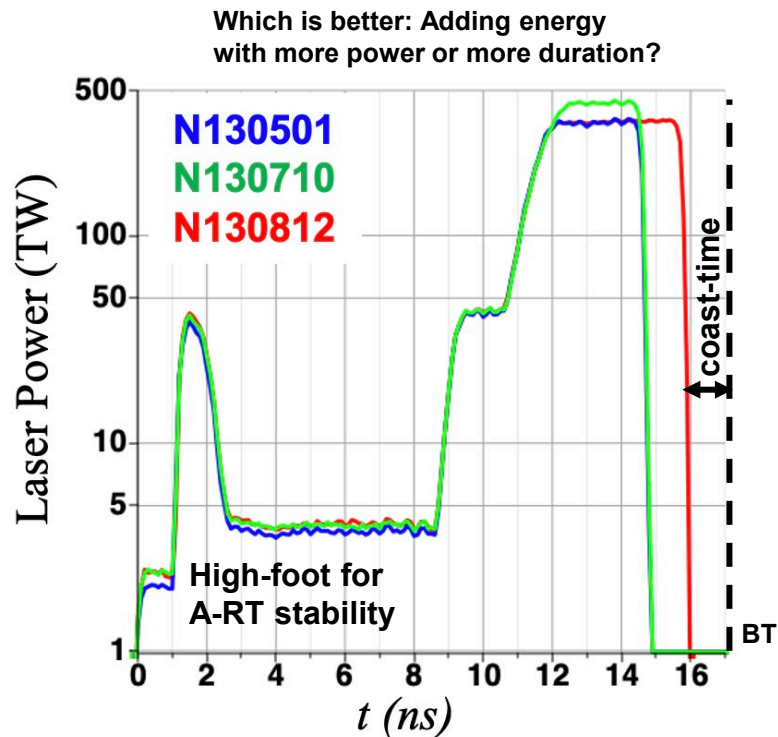


Fuel energy gain,  $G_{fuel}$ , is directly connected to the physics of the fusion plasma (Lawson Parameter)

$$\frac{Y_{total}}{E_{hs}} = \frac{5 \int_0^{\tau} m Q_a dt}{c_{DT} m T} \approx 4.6 \times 10^{26} p_{hs} \tau \frac{\langle \sigma v \rangle (T_{hs})}{T_{hs}^2}$$

with pressure in Gbar, T in keV, and  $\tau$  in s

# Even at the cost of peak power (at fixed energy), a longer duration of peak power (“lower coast time”) is beneficial



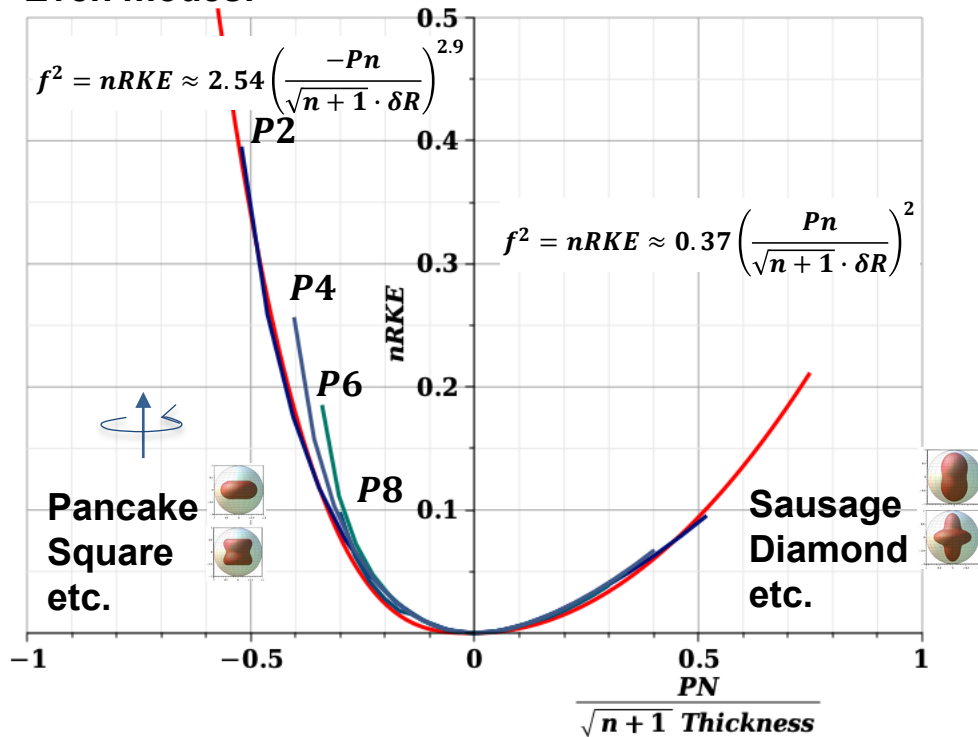
Optimal coast-time  $\ll$  hohlraum cooling time

Dittrich et al., PRL, 2014; Park et al., PRL, 2014; Hurricane, et al., Nature, 2014

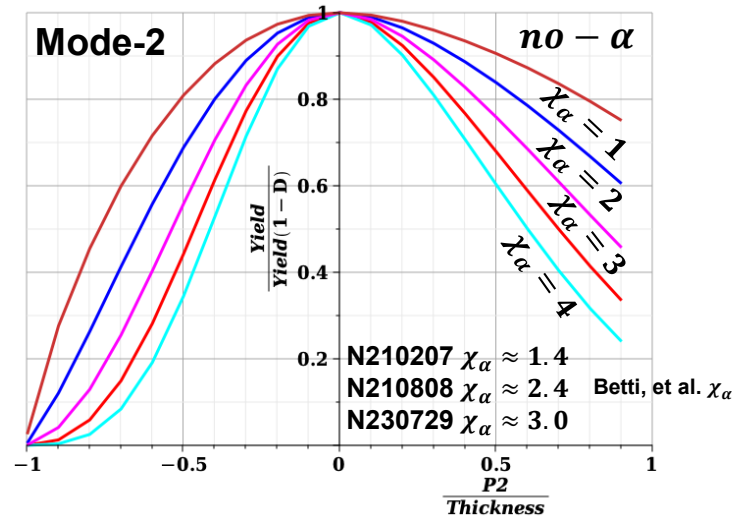
Effect on  $\rho R$  was previously noted: Zylstra, PoP (2014); Landen, PoP, (2012)

# Quantitatively, asymmetry sensitivity is *reduced* by shell thickness and *amplified* by alpha-particle self-heating

Even modes:



Asymmetry yield degradation is amplified by alpha-heating:

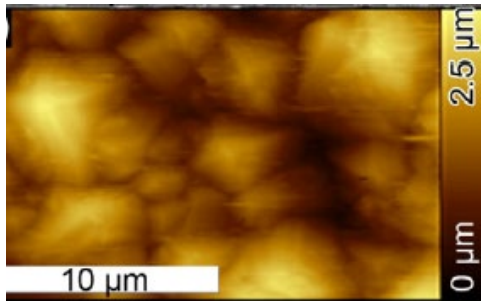


Asymmetry and shape observed in data: e.g. J. Ralph, et al, *Nat. Comm.*, (2024) 15:2975

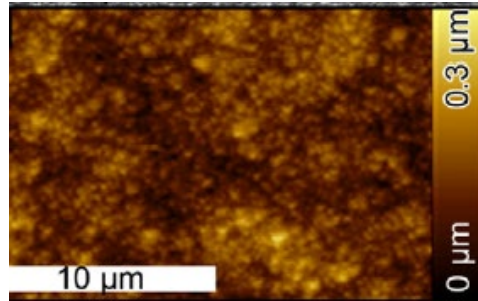
# Without a long-pulse driver capability, advanced light source facilities are likely limited to exploring target issues <10 Mbar

## Ablator structure

Micro-crystalline



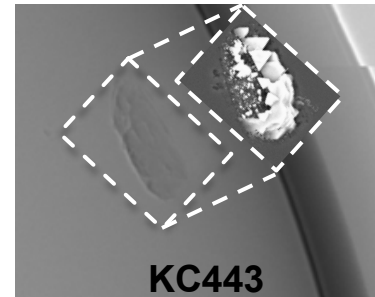
Nano-crystalline



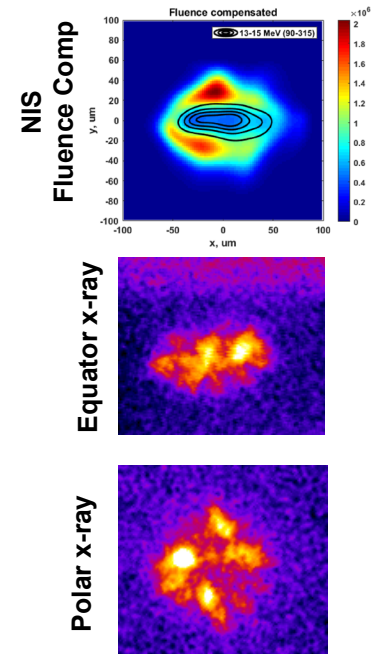
[C. Dawedeit et al. Diamond & Related Materials 40 75–81 (2013)]

## Defects

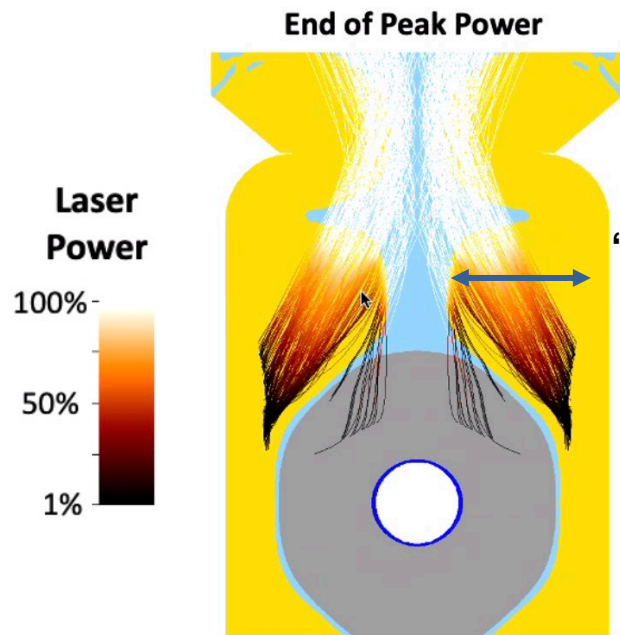
Inner surface defect



N181203  
KC468-01



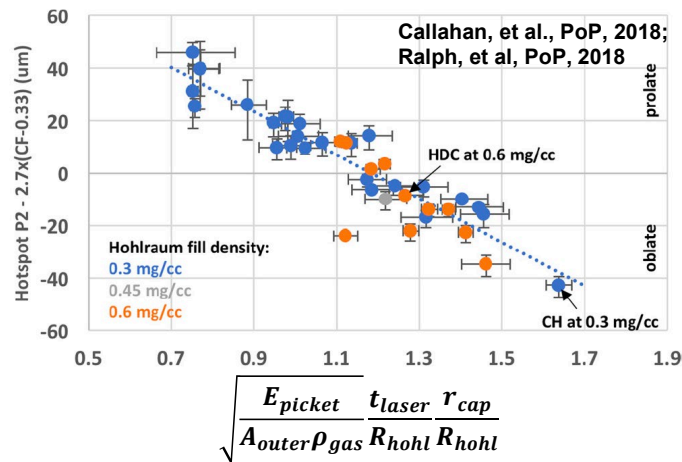
# Significantly improved understanding of the levers controlling IDD implosion symmetry obtained by 2018



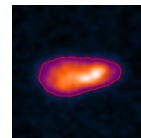
Legendre mode-2 (“P2”) empirical scaling:

$$x_{bubble} \sim \sqrt{\frac{E_{picket}}{\rho_{gas}}} t_{laser}$$

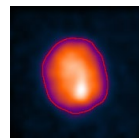
*Bubble blocks inner propagation*



$$\Delta\lambda = 0\text{\AA}$$



$$\Delta\lambda = 1\text{\AA}$$



A. L. Kritcher, et al *Phys. Rev. E* 98, 053206 (2018);  
L. Pickworth, et al, PoP (2020)